



SAPIENZA
UNIVERSITÀ DI ROMA

What can low energy quantum systems teach us about space and time?

Andrea Di Biagio

Advisors: Giovanni Montani and Carlo Rovelli

Tutor: Leonardo Gualtieri

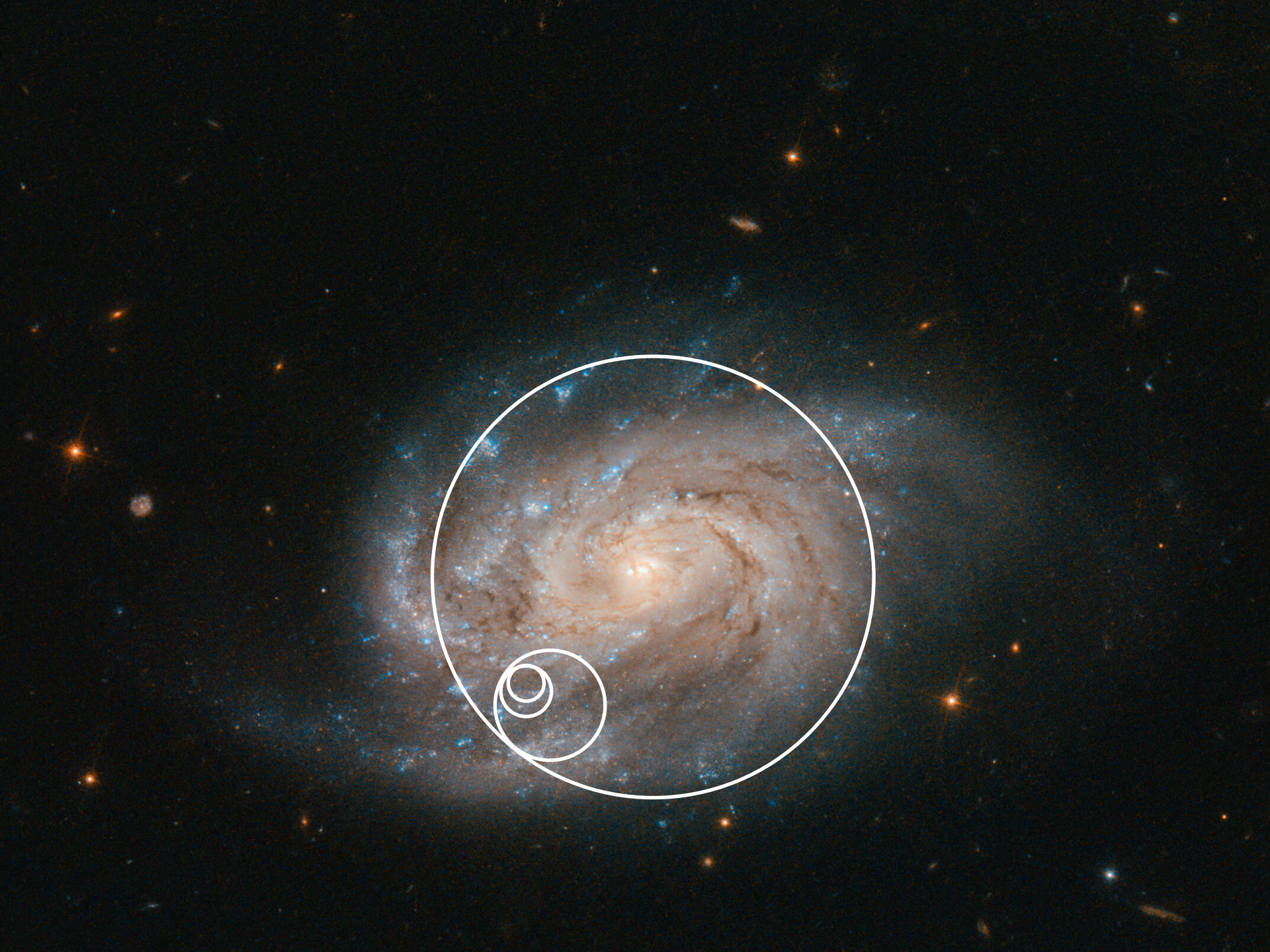
22 Feb 2022

PhD Viva 1

U. and Bhattacharjee, D. and Bhaumik, S. and Bilenko, I. A. and Billingsley, G. and Bini, S. and Birney, R. and Birnholz, O. and Bergers, S. and Bischli, M. and Biscoveanu, S. and Bisht, A. and ...

LIGO+Virgo

M. and Carrillo, G. and Carullo, G. and Carver, T. L. and Diaz, J. Casanueva and Casentini, C. and Castaldi, G. and Caudill, S. and Cavagli{\`a}, M. and Cavalieri, F. and Cavaliere, R. and ...



$$l_P = \sqrt{\frac{G\hbar}{c^3}} \approx 10^{-35} \text{ m}$$

$$E_P = \sqrt{\frac{\hbar c^5}{G}} \approx 10^{16} \text{ TeV}$$

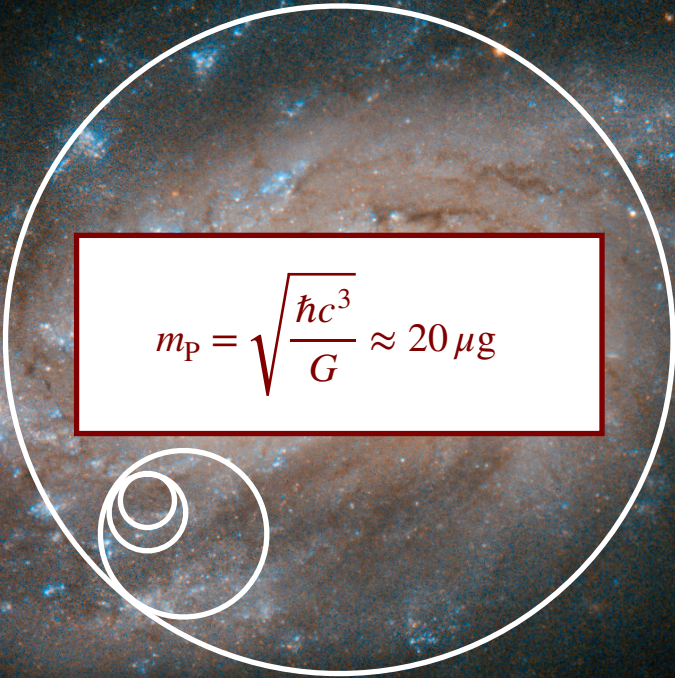
$$t_P = \sqrt{\frac{G\hbar}{c^5}} \approx 10^{-44} \text{ s}$$



$$l_P = \sqrt{\frac{G\hbar}{c^3}} \approx 10^{-35} \text{ m}$$

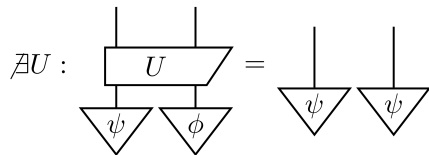
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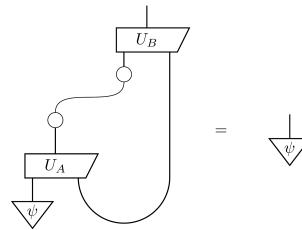

$$m_P = \sqrt{\frac{\hbar c^3}{G}} \approx 20 \mu\text{g}$$

Quantum Information and Foundations

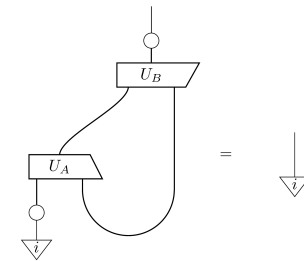
QI: Information-processing capabilities afforded by quantum systems



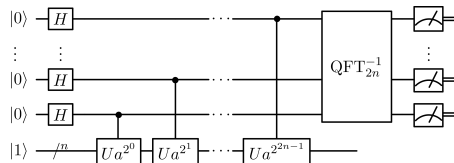
No-cloning theorem



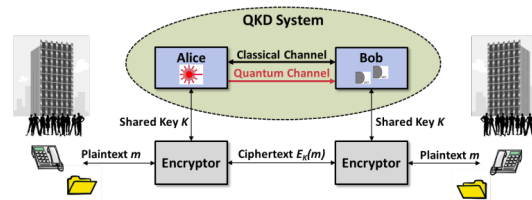
Quantum Teleportation



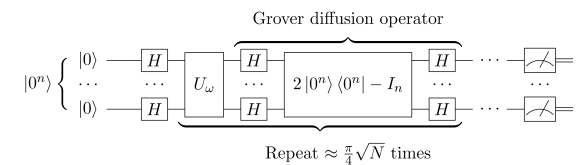
Superdense coding



Shor's algorithm



Quantum key distribution



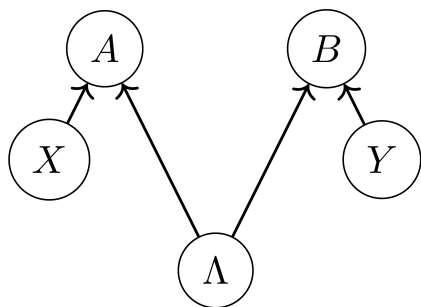
Grover's algorithm

Quantum Information and Foundations

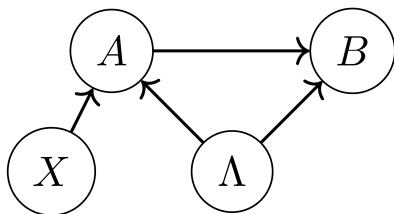
QF: study of the counterintuitive properties of QM

Quantum Information and Foundations

QF: study of the counterintuitive properties of QM

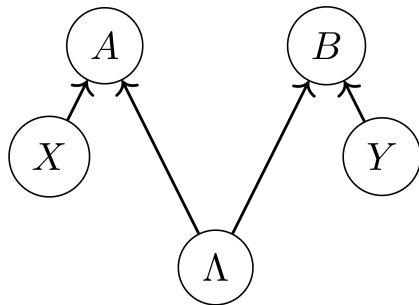


Violation of causal inequalities



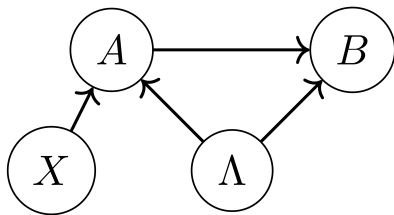
Quantum Information and Foundations

QF: study of the counterintuitive properties of QM

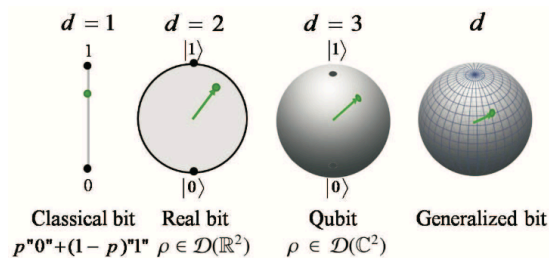


1. **Finiteness.** If a system carries one bit of information, then each state is characterised by the outcome probabilities of a finite set of measurements.
2. **Local tomography.** The state of a composite system is fully characterised by the statistics of measurements performed on the subsystems.
3. **Equivalence of subspaces.** Systems that carry the same amount of information have isomorphic state spaces.
4. **Symmetry.** Any pure state can be reversibly transformed into any other pure state.
5. **All measurements are allowed.** Every mathematically well defined effect on a system carrying one bit corresponds to a possible measurement.

Violation of causal inequalities

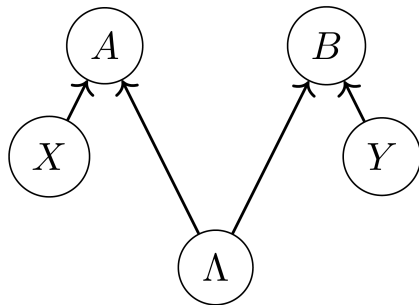


Reconstructions of quantum theory from physical principles

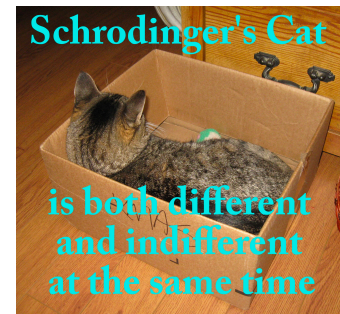


Quantum Information and Foundations

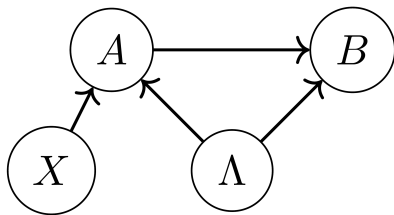
QF: study of the counterintuitive properties of QM



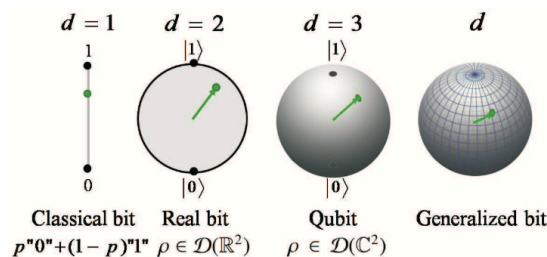
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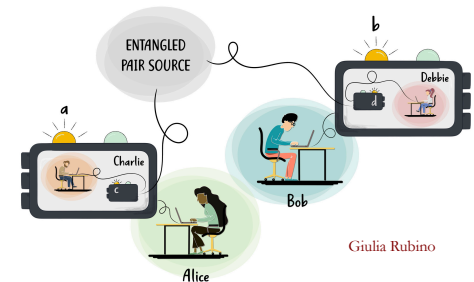
Violation of causal inequalities



Reconstructions of quantum theory from physical principles



Interpretations of quantum theory



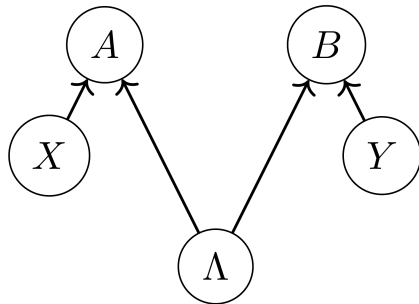
Giulia Rubino

Quantum Information and Foundations

Significant-Loophole-Free Test of Bell's Theorem with Entangled Photons

Marissa Giustina, Marijn A. M. Versteegh, Sören Wengerowsky, Johannes Handsteiner, Armin Hochrainer, Kevin Phelan, Fabian Steinlechner, Johannes Kofler, Jan-Åke Larsson, Carlos Abellán, Waldimar Amaya, Valerio Pruneri, Morgan W. Mitchell, Jörn Beyer, Thomas Gerrits, Adriana E. Lita, Lynden K. Shalm, Sae Woo Nam, Thomas Scheidl, Rupert Ursin, Bernhard Wittmann, and Anton Zeilinger
 Phys. Rev. Lett. **115**, 250401 – Published 16 December 2015

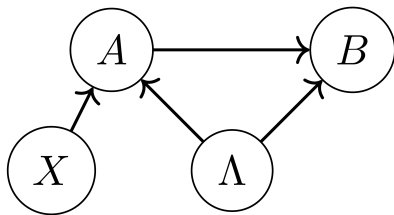
of QM



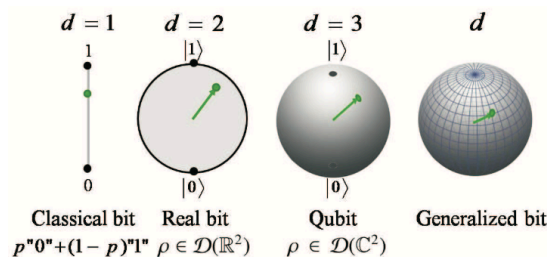
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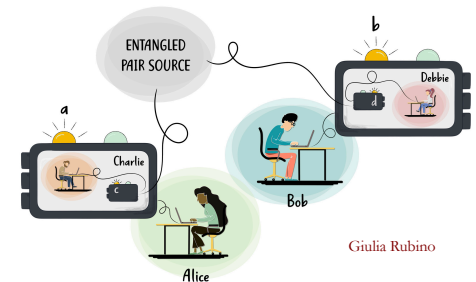
Violation of causal inequalities



Reconstructions of quantum theory from physical principles



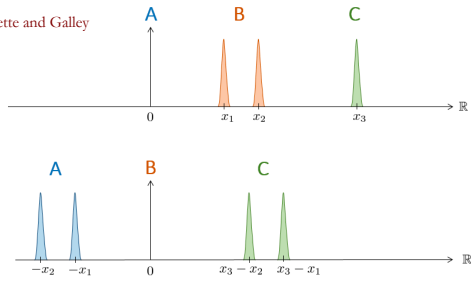
Interpretations of quantum theory



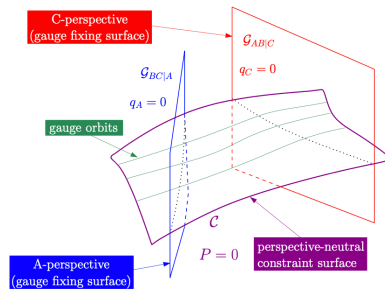
QI and Spacetime

QI and Spacetime

de la Hamette and Galley

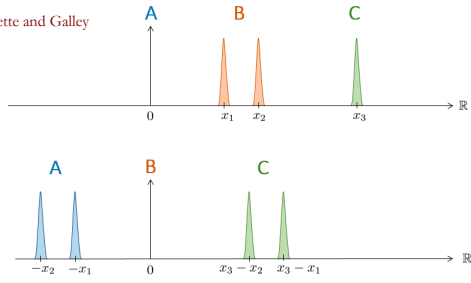


Quantum Reference Frames

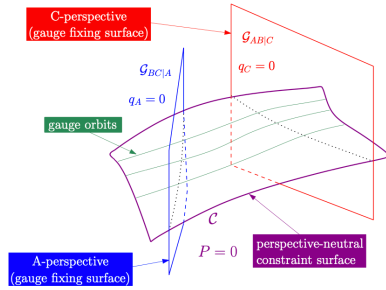


QI and Spacetime

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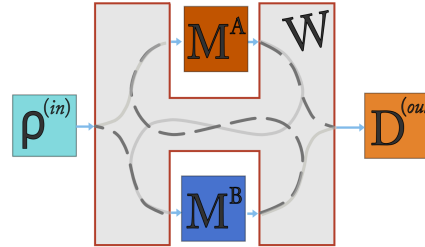


Quantum Reference Frames

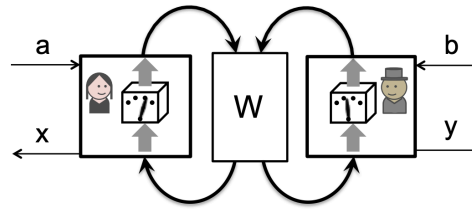


Vanrietvelde, Höhn, Giacomini, Castro-Ruiz

Rubino et.al.



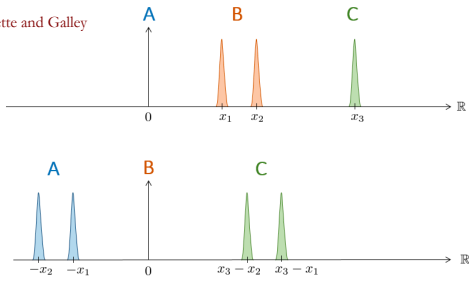
Indefinite causality



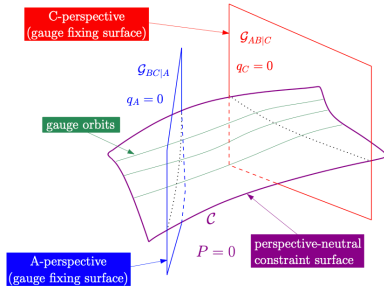
Oreshkov, Costa, Brukner

QI and Spacetime

de la Hamette and Galley

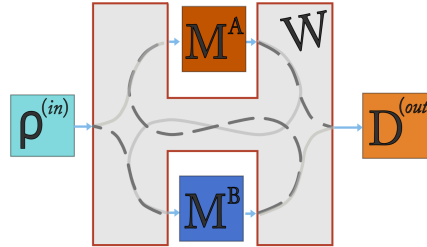


Quantum Reference Frames

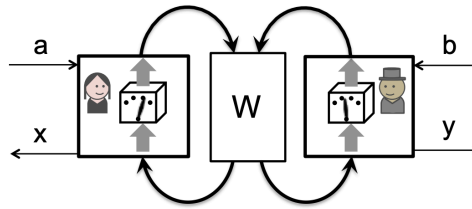


Vannietvelde, Höhn, Giacomini, Castro-Ruiz

Rubino et.al.

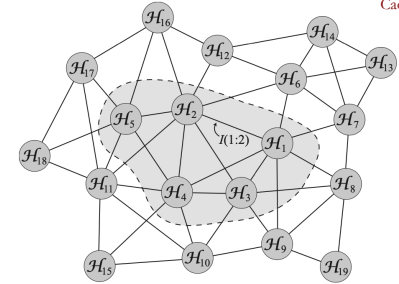


Indefinite causality

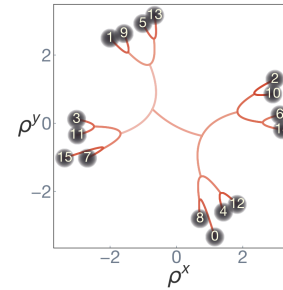


Oreshkov, Costa, Brukner

Cao, Carrol, Michalakis



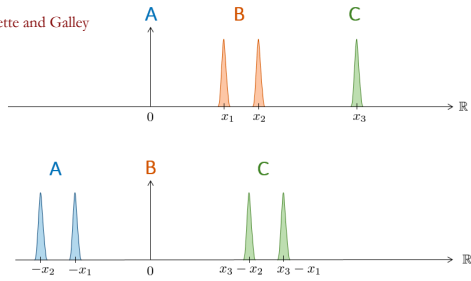
Geometry from entanglement



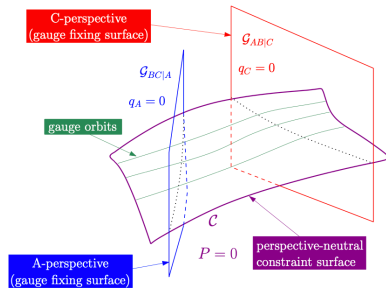
Periwal et.al.

QI and Spacetime

de la Hamette and Galley

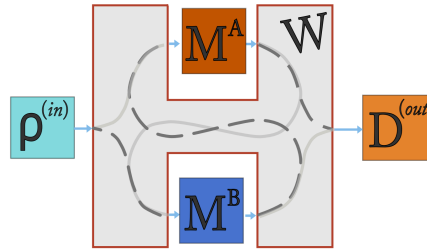


Quantum Reference Frames

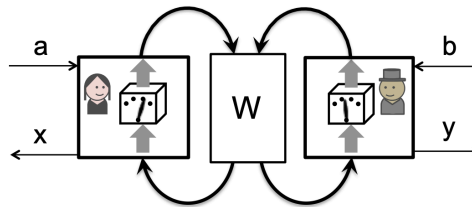


Vannietvelde, Höhn, Giacomini, Castro-Ruiz

Rubino et.al.

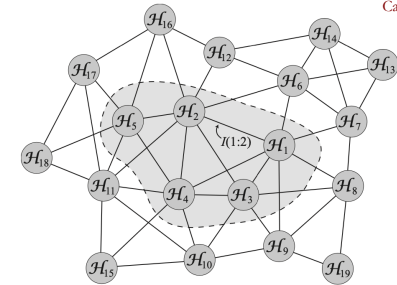


Indefinite causality

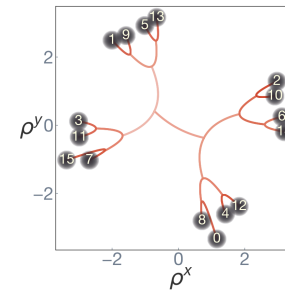


Oreshkov, Costa, Brukner

Cao, Carrol, Michalakis



Geometry from entanglement



Periwal et.al.

Low energy tests of quantum gravity!

Krisnanda et.al.

- **Part I: Quantum gravity (and beyond) in the lab**
- **Part II: Conceptual investigations**

- **Part I: Quantum gravity (and beyond) in the lab**
 - **Gravitationally Mediated Entanglement (GME)**
 - **Optical Simulation of a GME experiment**
 - **Computing the phases from first principles**
 - **An experiment to test the discreteness of time**

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Gravitationally Induced Entanglement between Two Massive Particles is Sufficient Evidence of Quantum Effects in Gravity

C. Marletto and V. Vedral
Phys. Rev. Lett. **119**, 240402 – Published 13 December 2017

Spin Entanglement Witness for Quantum Gravity

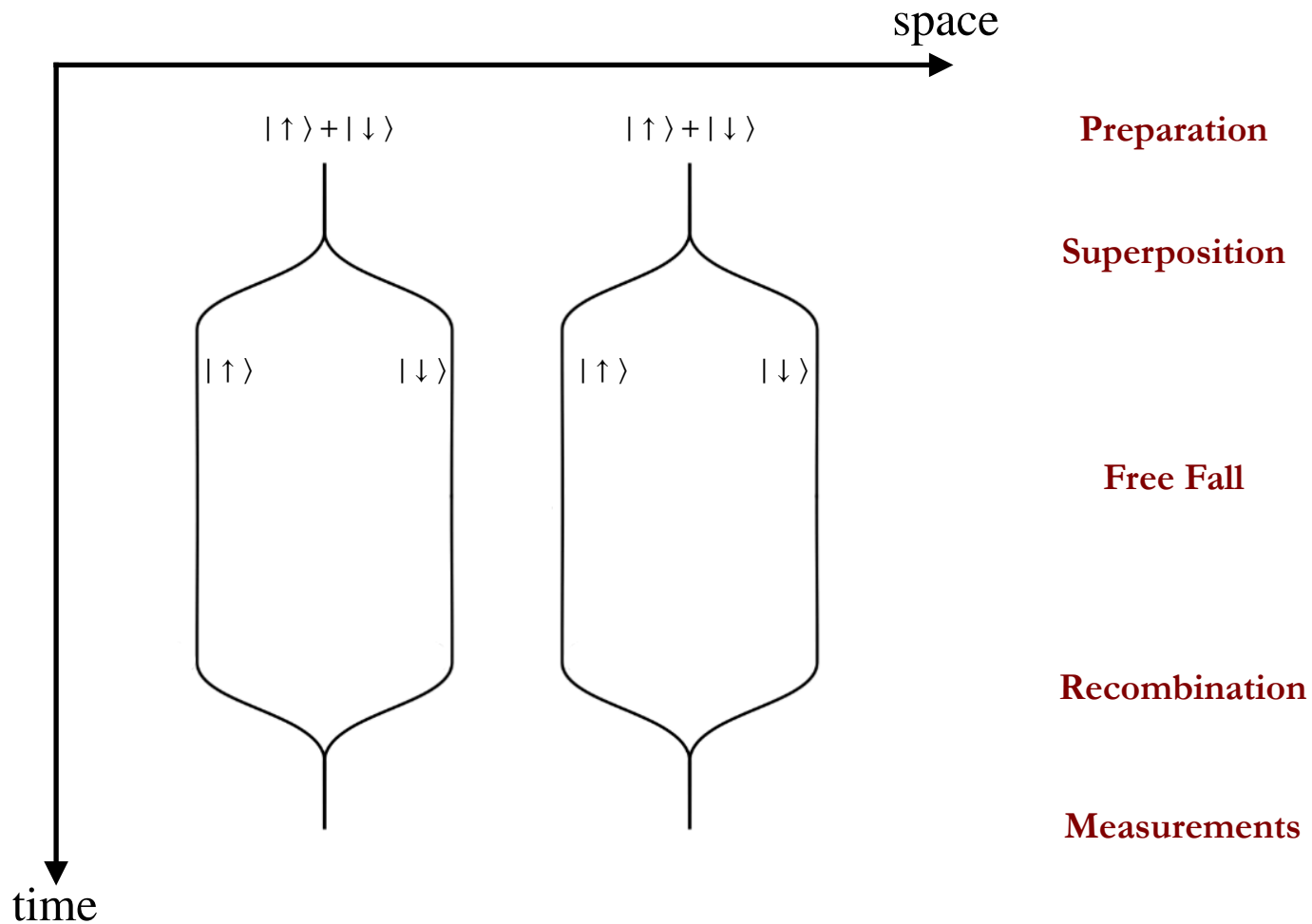
Sougato Bose, Anupam Mazumdar, Gavin W. Morley, Hendrik Ulbricht, Marko Toroš, Mauro Paternostro, Andrew A. Geraci, Peter F. Barker, M. S. Kim, and Gerard Milburn
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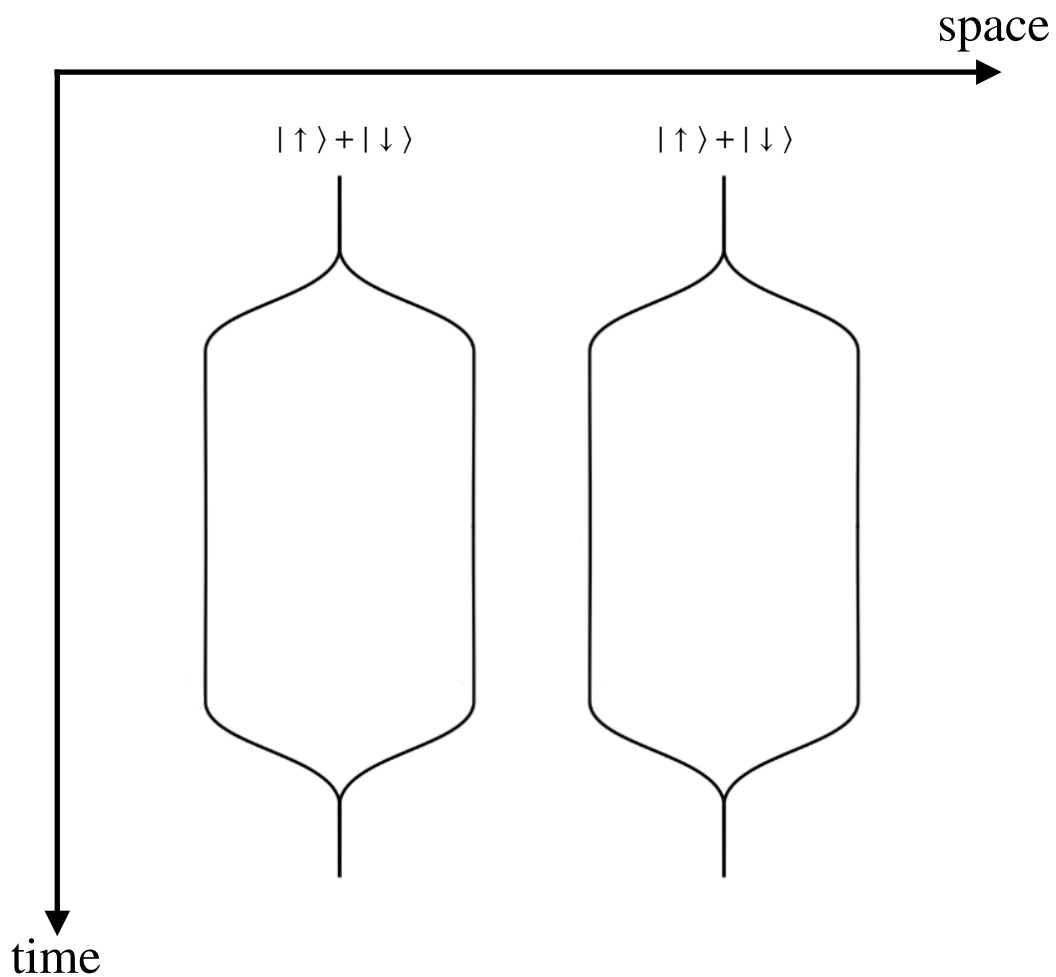
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 Phys. Rev. Lett. **119**, 240401 – Published 13 December 2017



GME



Preparation

$$(|\uparrow\uparrow\rangle + |\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle + |\downarrow\downarrow\rangle) |g_{cc}\rangle$$

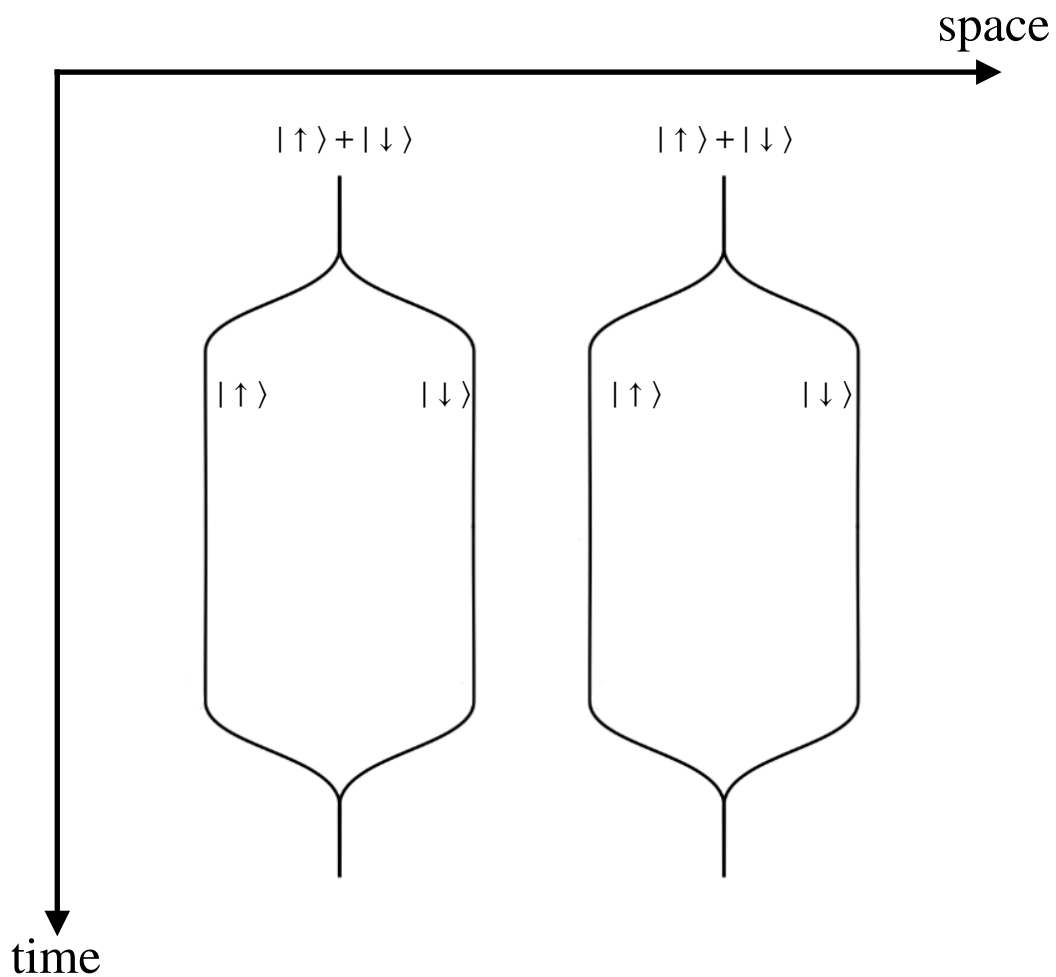
Superposition

Free Fall

Recombination

Measurements

GME



Preparation

$$(|\uparrow\uparrow\rangle + |\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle + |\downarrow\downarrow\rangle) |g_{cc}\rangle$$

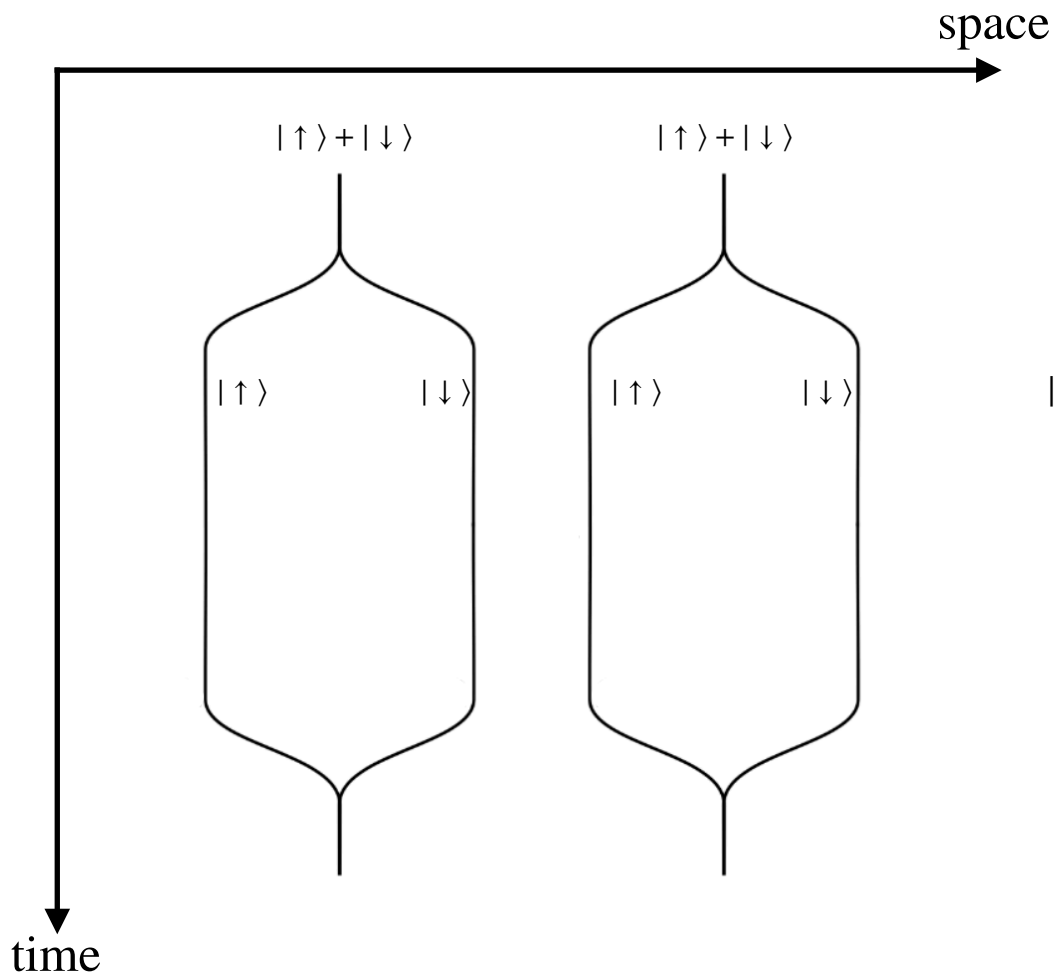
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Superposition

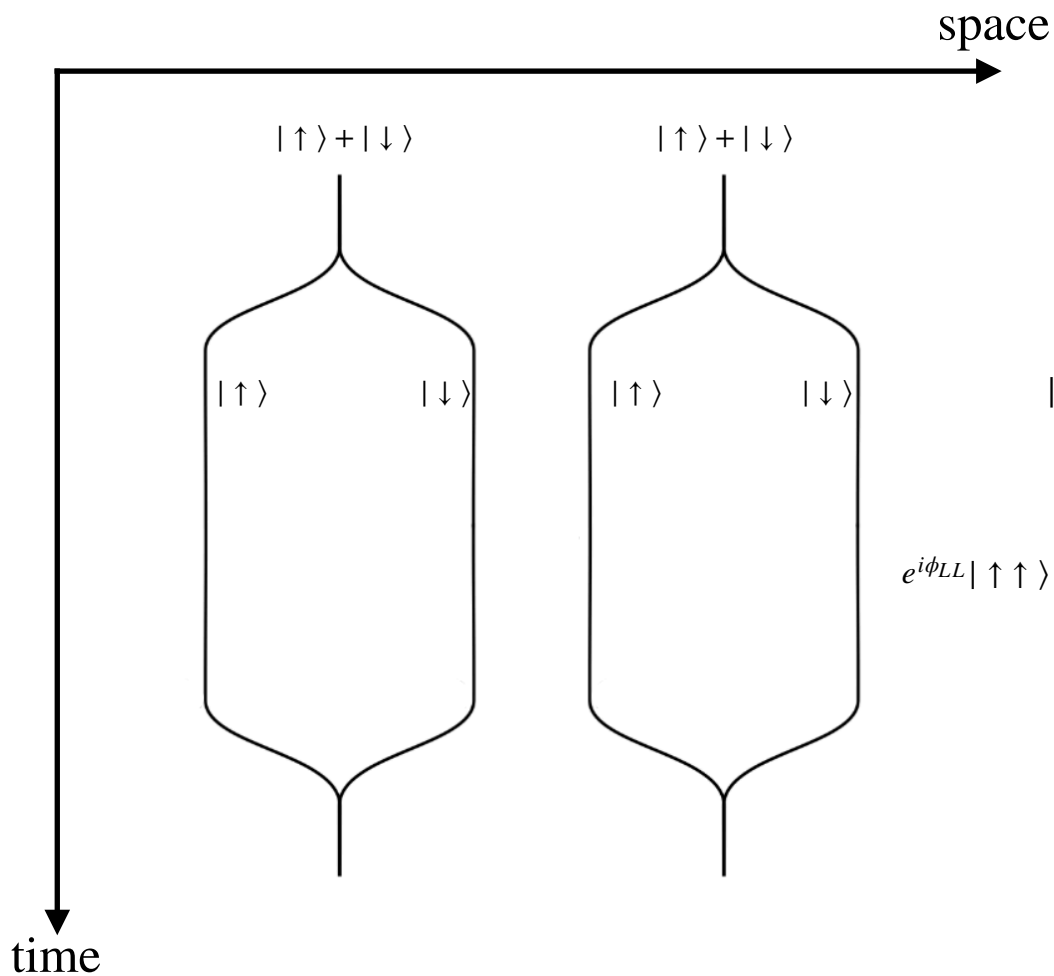
$$|\uparrow\uparrow\rangle |g_{LL}\rangle + |\uparrow\downarrow\rangle |g_{LR}\rangle + |\downarrow\uparrow\rangle |g_{RL}\rangle + |\downarrow\downarrow\rangle |g_{RR}\rangle$$

Free Fall

Recombination

Measurements

GME



Preparation

$$(|\uparrow\uparrow\rangle + |\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle + |\downarrow\downarrow\rangle) |g_{CC}\rangle$$

Superposition

$$|\uparrow\uparrow\rangle |g_{LL}\rangle + |\uparrow\downarrow\rangle |g_{LR}\rangle + |\downarrow\uparrow\rangle |g_{RL}\rangle + |\downarrow\downarrow\rangle |g_{RR}\rangle$$

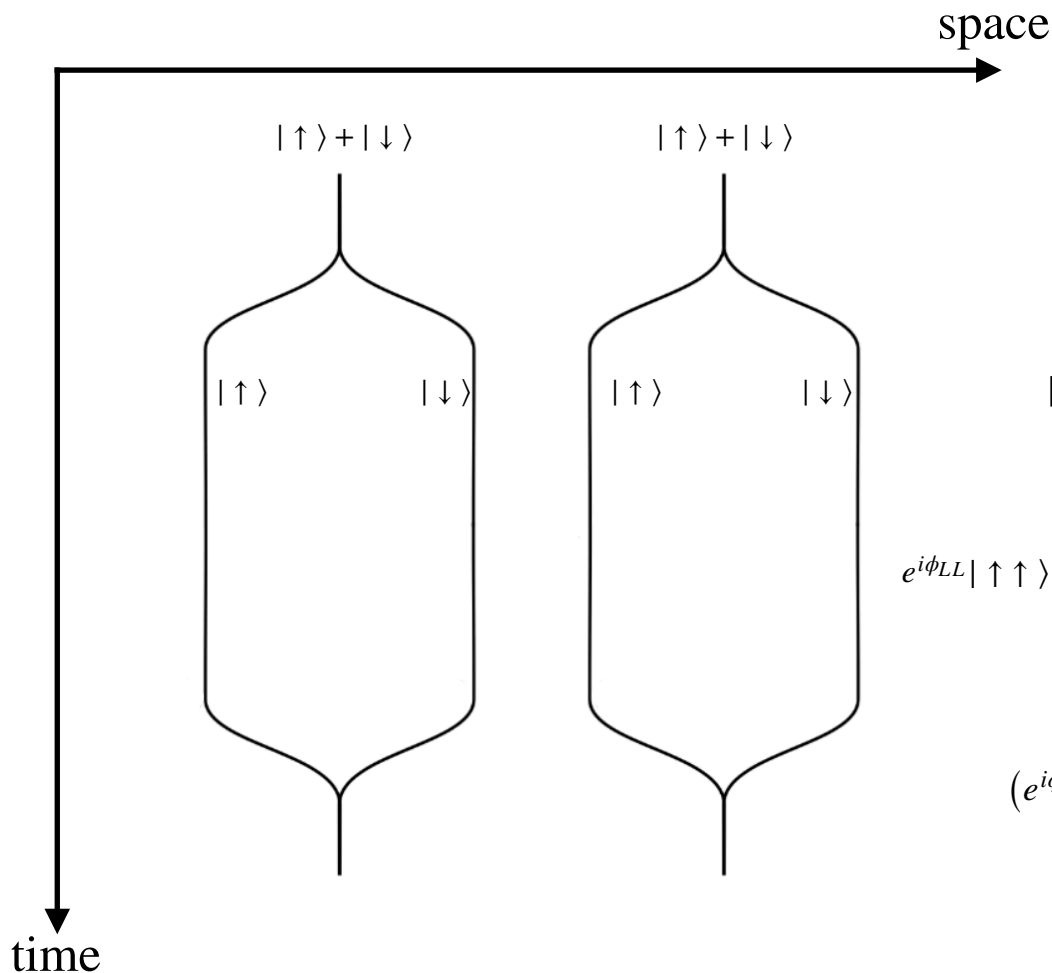
Free Fall

$$e^{i\phi_{LL}} |\uparrow\uparrow\rangle |g_{LL}\rangle + e^{i\phi_{LR}} |\uparrow\downarrow\rangle |g_{LR}\rangle + e^{i\phi_{RL}} |\downarrow\uparrow\rangle |g_{RL}\rangle + e^{i\phi_{RR}} |\downarrow\downarrow\rangle |g_{RR}\rangle$$

Recombination

Measurements

GME



Preparation

$$(|\uparrow\uparrow\rangle + |\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle + |\downarrow\downarrow\rangle) |g_{CC}\rangle$$

Superposition

$$|\uparrow\uparrow\rangle |g_{LL}\rangle + |\uparrow\downarrow\rangle |g_{LR}\rangle + |\downarrow\uparrow\rangle |g_{RL}\rangle + |\downarrow\downarrow\rangle |g_{RR}\rangle$$

Free Fall

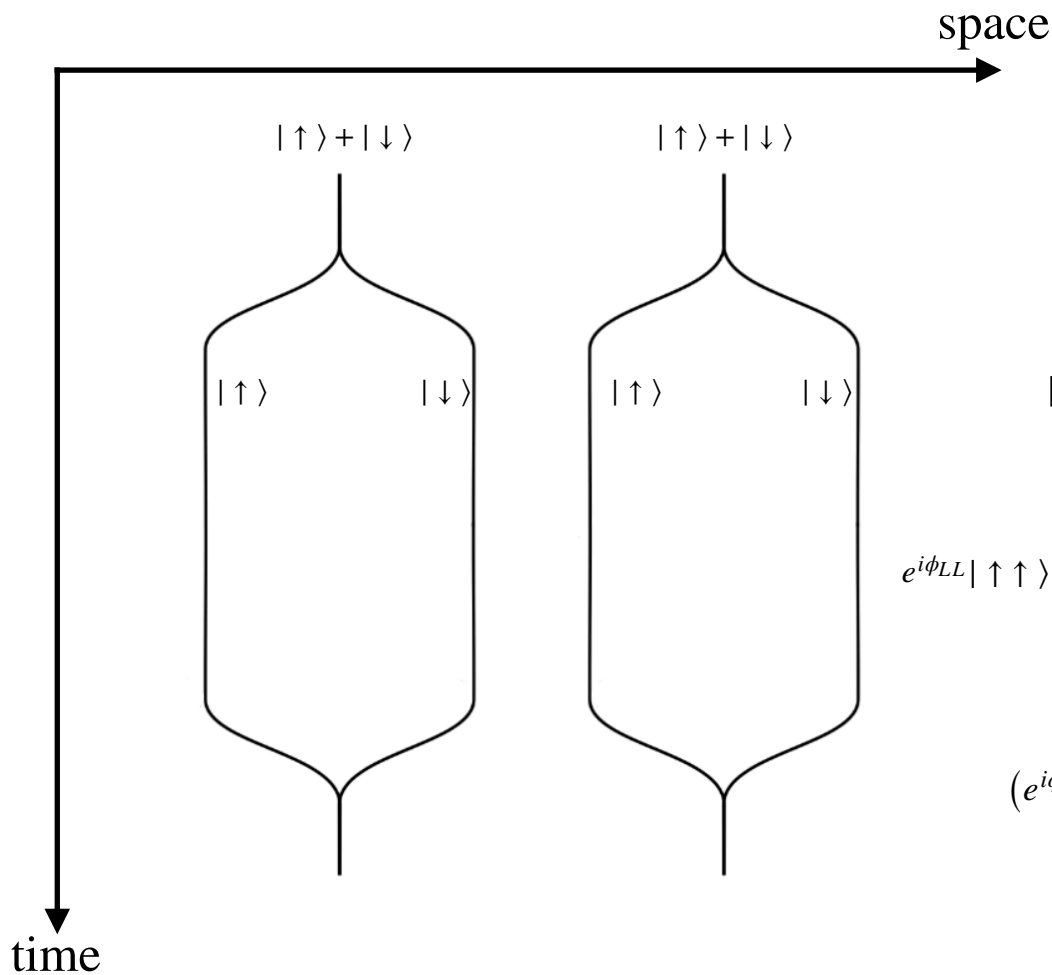
$$e^{i\phi_{LL}} |\uparrow\uparrow\rangle |g_{LL}\rangle + e^{i\phi_{LR}} |\uparrow\downarrow\rangle |g_{LR}\rangle + e^{i\phi_{RL}} |\downarrow\uparrow\rangle |g_{RL}\rangle + e^{i\phi_{RR}} |\downarrow\downarrow\rangle |g_{RR}\rangle$$

Recombination

$$(e^{i\phi_{LL}} |\uparrow\uparrow\rangle + e^{i\phi_{LR}} |\uparrow\downarrow\rangle + e^{i\phi_{RL}} |\downarrow\uparrow\rangle + e^{i\phi_{RR}} |\downarrow\downarrow\rangle) |g_{CC}\rangle$$

Measurements

GME



Preparation

$$(|\uparrow\uparrow\rangle + |\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle + |\downarrow\downarrow\rangle) |g_{CC}\rangle$$

Superposition

$$|\uparrow\uparrow\rangle |g_{LL}\rangle + |\uparrow\downarrow\rangle |g_{LR}\rangle + |\downarrow\uparrow\rangle |g_{RL}\rangle + |\downarrow\downarrow\rangle |g_{RR}\rangle$$

Free Fall

$$e^{i\phi_{LL}} |\uparrow\uparrow\rangle |g_{LL}\rangle + e^{i\phi_{LR}} |\uparrow\downarrow\rangle |g_{LR}\rangle + e^{i\phi_{RL}} |\downarrow\uparrow\rangle |g_{RL}\rangle + e^{i\phi_{RR}} |\downarrow\downarrow\rangle |g_{RR}\rangle$$

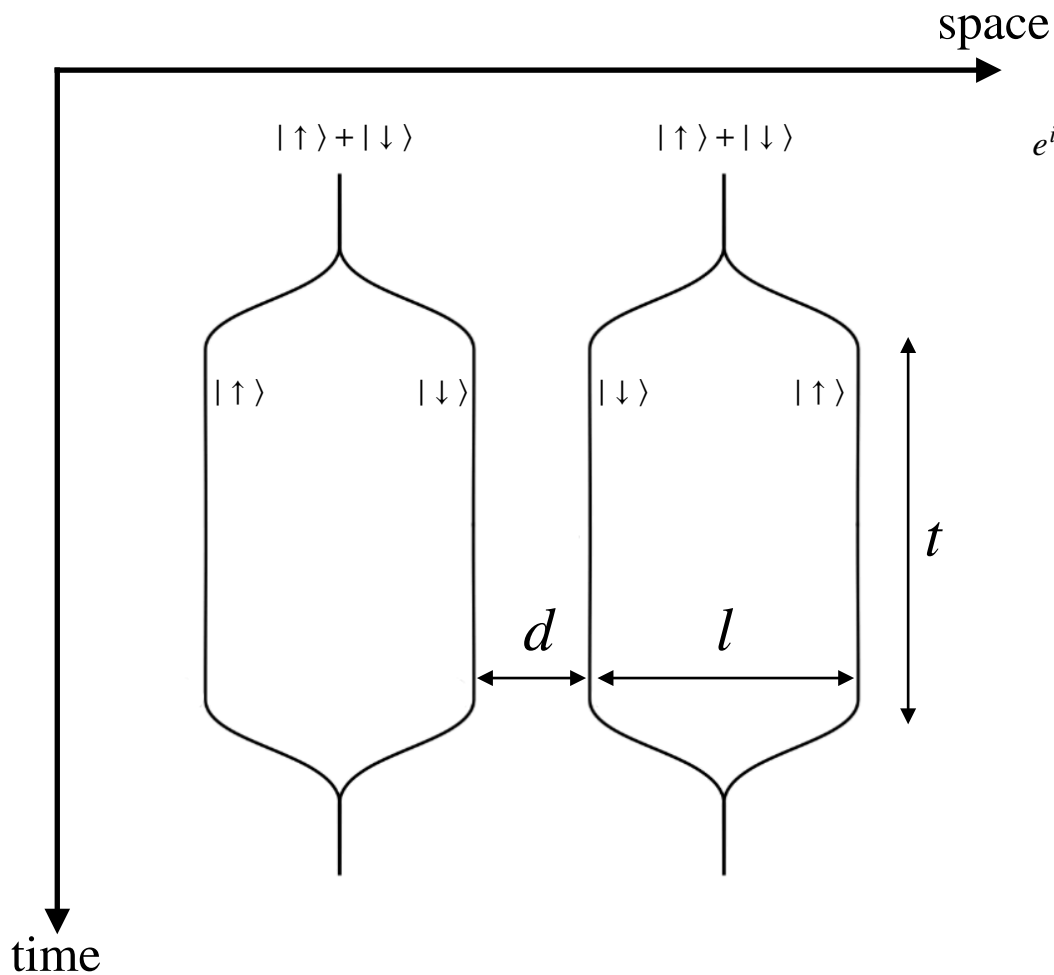
Recombination

$$(e^{i\phi_{LL}} |\uparrow\uparrow\rangle + e^{i\phi_{LR}} |\uparrow\downarrow\rangle + e^{i\phi_{RL}} |\downarrow\uparrow\rangle + e^{i\phi_{RR}} |\downarrow\downarrow\rangle) |g_{CC}\rangle$$

Measurements

$$e^{i\phi_{LL}} |\uparrow\uparrow\rangle + e^{i\phi_{LR}} |\uparrow\downarrow\rangle + e^{i\phi_{RL}} |\downarrow\uparrow\rangle + e^{i\phi_{RR}} |\downarrow\downarrow\rangle$$

GME

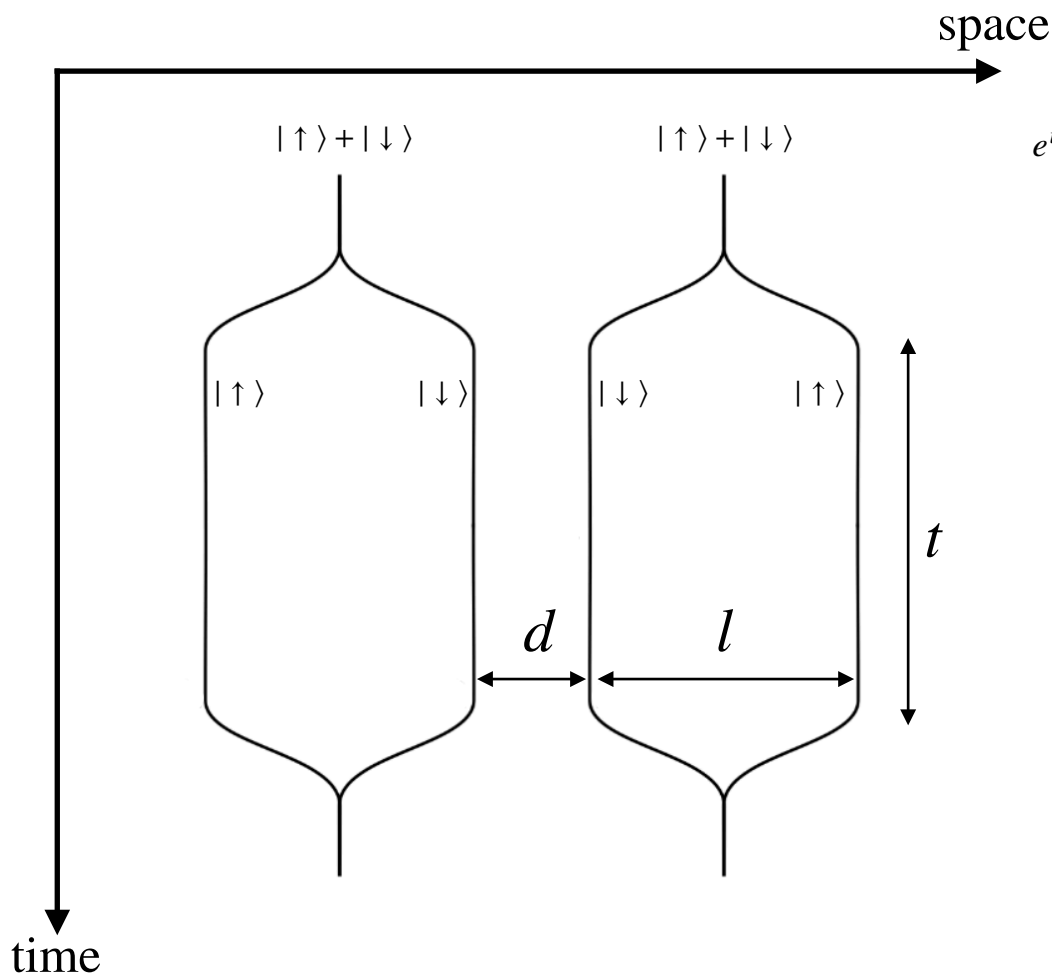


$$e^{i\phi_{LR}}|\uparrow\uparrow\rangle + e^{i\phi_{LL}}|\uparrow\downarrow\rangle + e^{i\phi_{RR}}|\downarrow\uparrow\rangle + e^{i\phi_{RL}}|\downarrow\downarrow\rangle$$

$$\phi_{LR} = \frac{Gm^2}{d+2l} \frac{t}{\hbar} \quad \phi_{RR} = \frac{Gm^2}{d+l} \frac{t}{\hbar} = \phi_{LL}$$

$$\phi_{RL} = \frac{Gm^2}{d} \frac{t}{\hbar}$$

GME



$$e^{i\phi_{LR}}|\uparrow\uparrow\rangle + e^{i\phi_{LL}}|\uparrow\downarrow\rangle + e^{i\phi_{RR}}|\downarrow\uparrow\rangle + e^{i\phi_{RL}}|\downarrow\downarrow\rangle$$

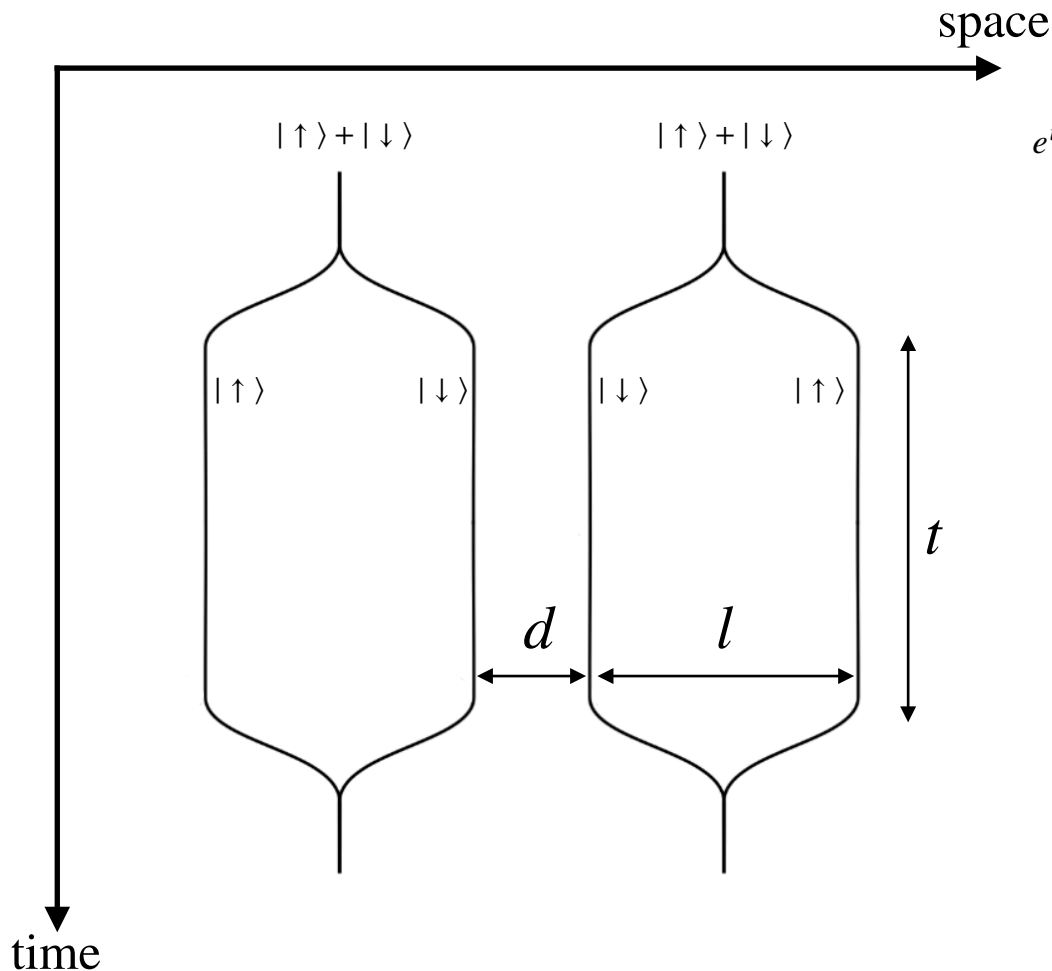
$$\phi_{LR} = \frac{Gm^2 t}{d + 2l \hbar} \quad \phi_{RR} = \frac{Gm^2 t}{d + l \hbar} = \phi_{LL}$$

$$\phi_{RL} = \frac{Gm^2 t}{d \hbar}$$

$$m \approx 10^{-14} \text{kg} \approx 10^{-6} m_p$$

$$t \approx 1 \text{s} \quad d \approx 200 \mu\text{m} \quad l \approx 250 \mu\text{m}$$

$$\Rightarrow \Delta\phi = \phi_{RL} + \phi_{LR} - 2\phi_{LL} \approx 1$$



$$e^{i\phi_{LR}}|\uparrow\uparrow\rangle + e^{i\phi_{LL}}|\uparrow\downarrow\rangle + e^{i\phi_{RR}}|\downarrow\uparrow\rangle + e^{i\phi_{RL}}|\downarrow\downarrow\rangle$$

$$\phi_{LR} = \frac{Gm^2}{d+2l} \frac{t}{\hbar} \quad \phi_{RR} = \frac{Gm^2}{d+l} \frac{t}{\hbar} = \phi_{LL}$$

$$\phi_{RL} = \frac{Gm^2}{d} \frac{t}{\hbar}$$

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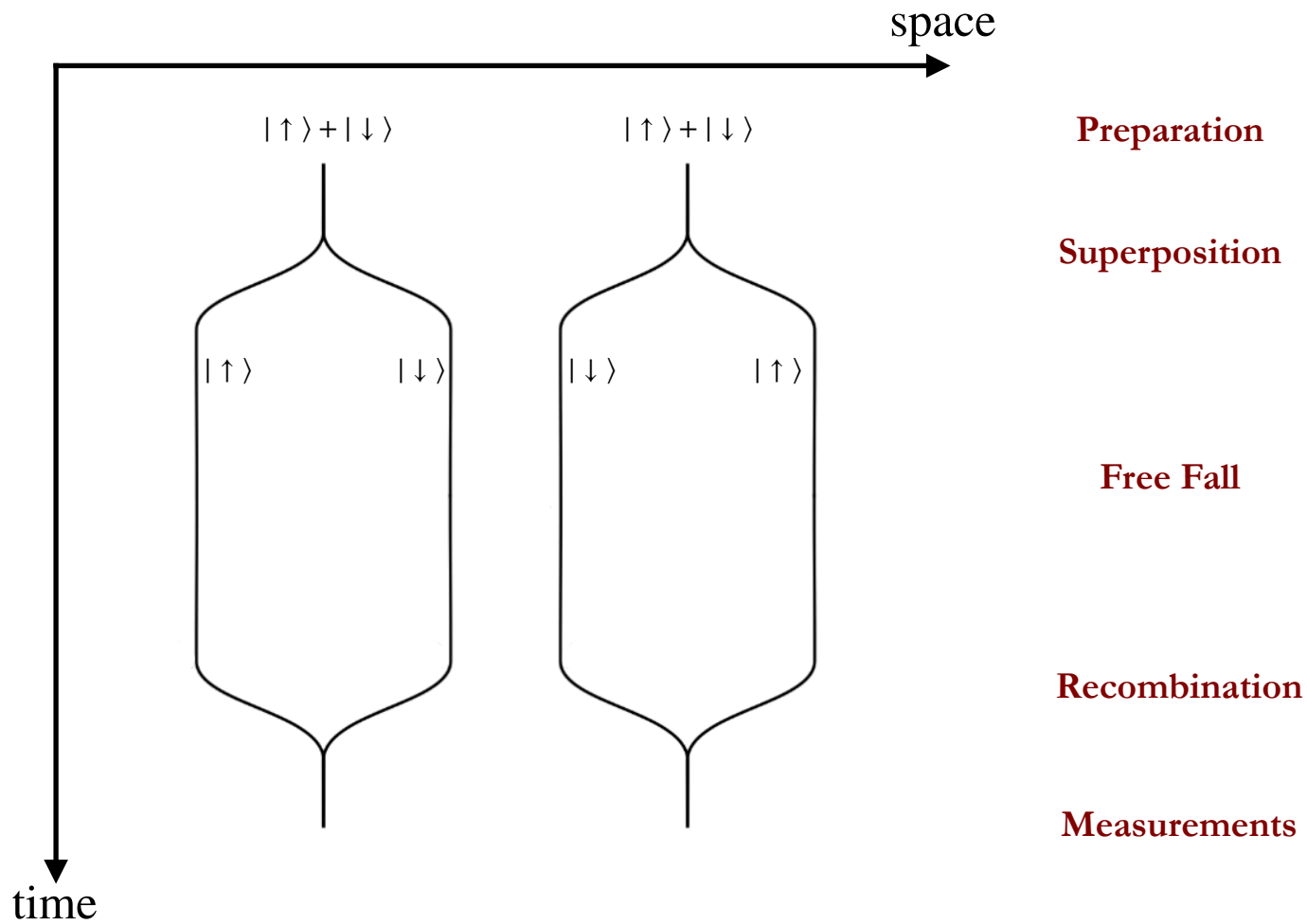
A prediction of linearised quantum gravity.

Not explainable in terms of semi-classical gravity

$$G_{\mu\nu} \propto \langle \hat{T}_{\mu\nu} \rangle$$

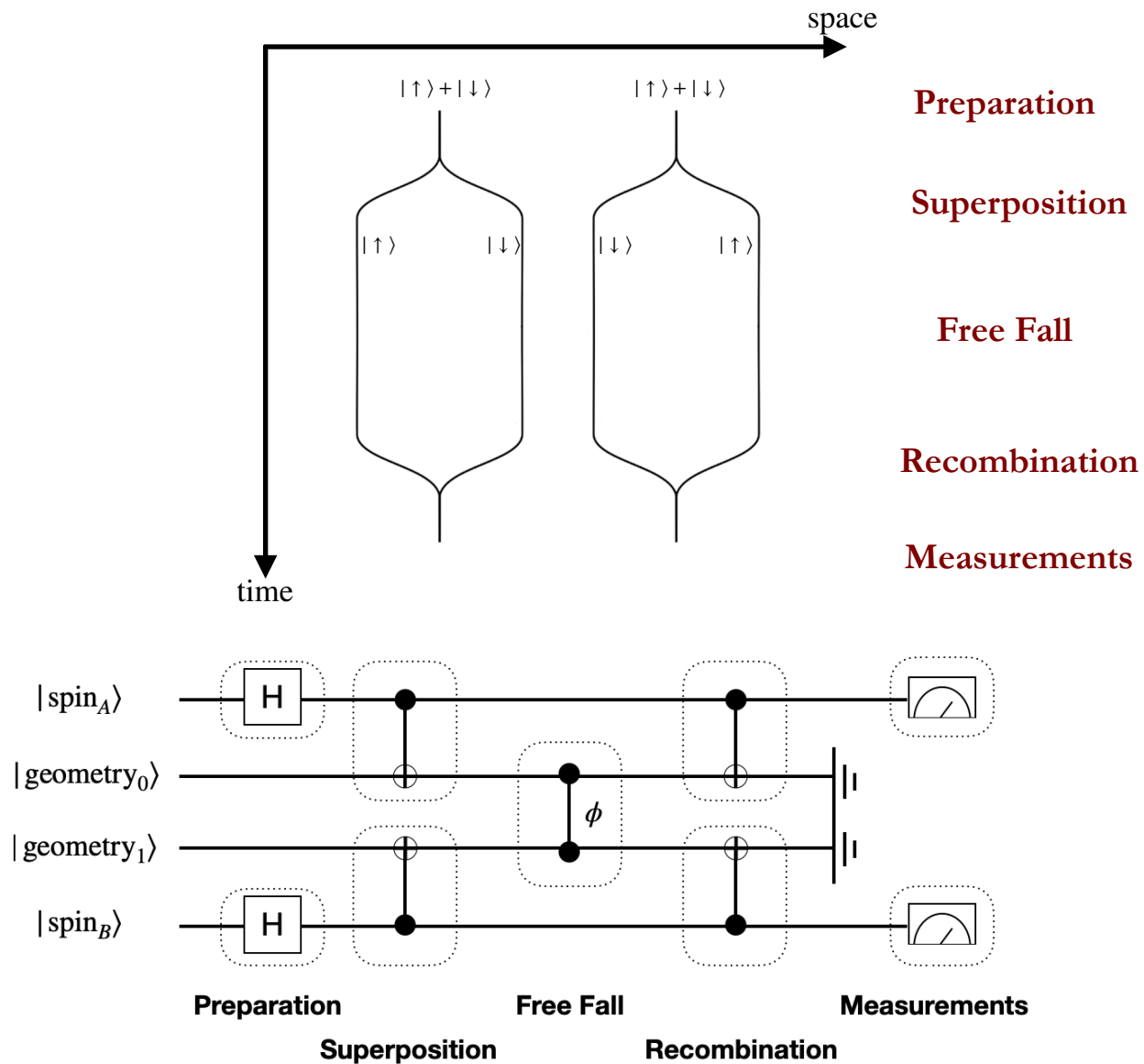
- Part I: Quantum gravity (and beyond) in the lab
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Simulating GME

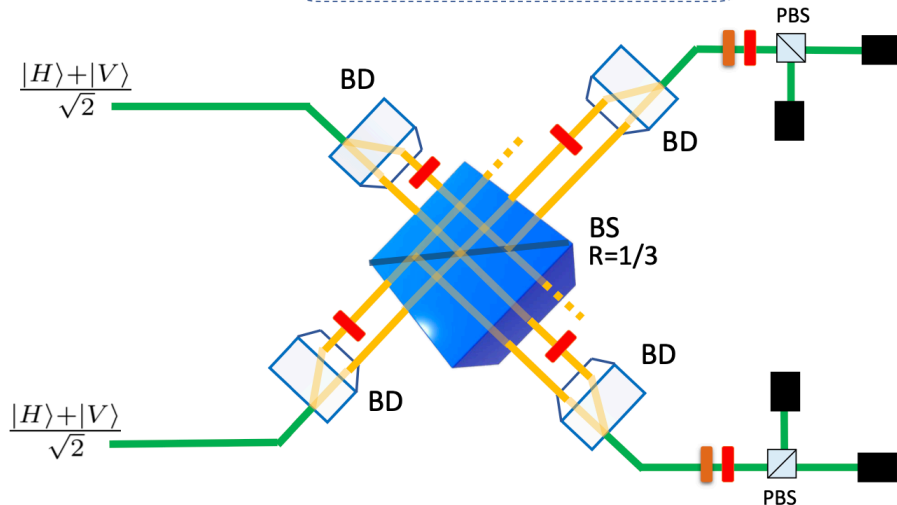
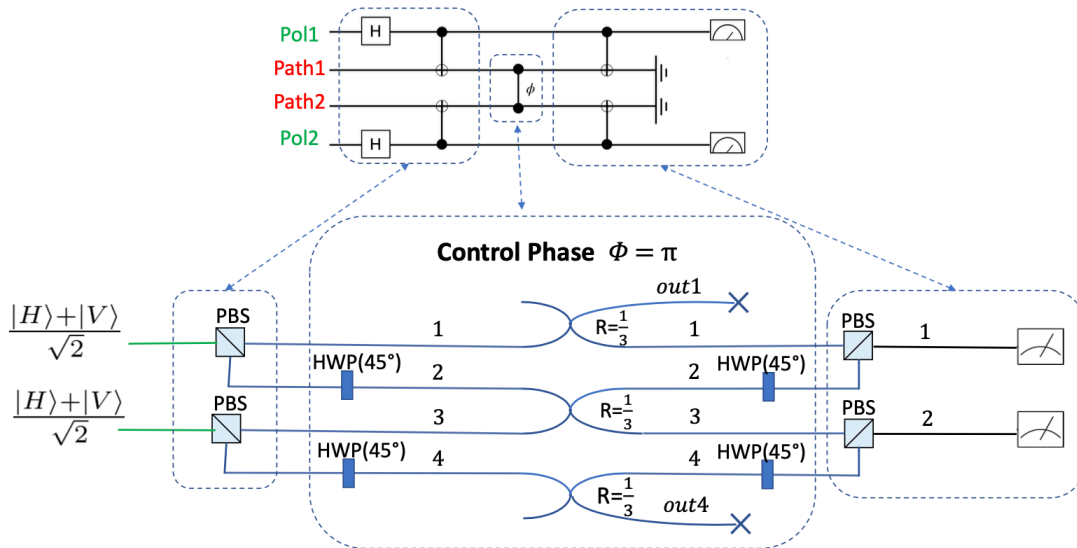


$$\mathbb{C}^2 \otimes \mathbb{C}^4 \otimes \mathbb{C}^2 = \mathcal{H}_{\text{spin}_A} \otimes \mathcal{H}_{\text{geometry}} \otimes \mathcal{H}_{\text{spin}_B}$$

Quantum Circuit



Optics Simulation



Two-photon scheme:

spins \rightarrow polarisation

geometry \rightarrow path

Entanglement witness

separable $\implies W \geq 0$

$$W^{\text{exp}} = -0.514(2)$$

CHSH inequality violation

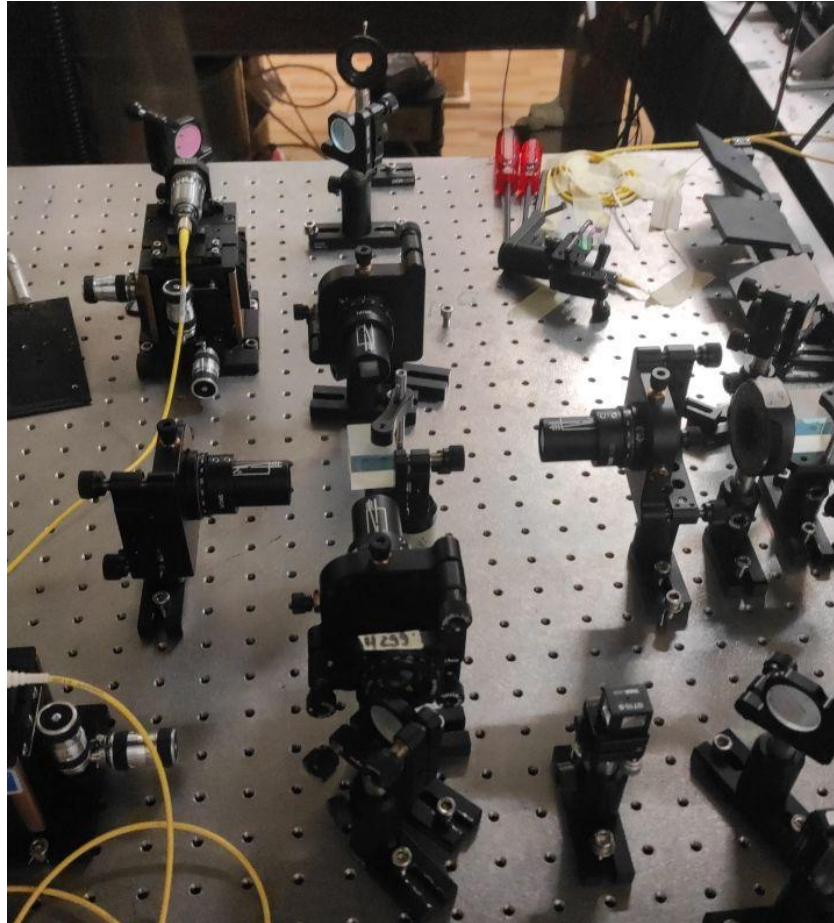
classical bound $S \leq 2$

$$S^{\text{exp}} = 2.401(15)$$

Optics Simulation

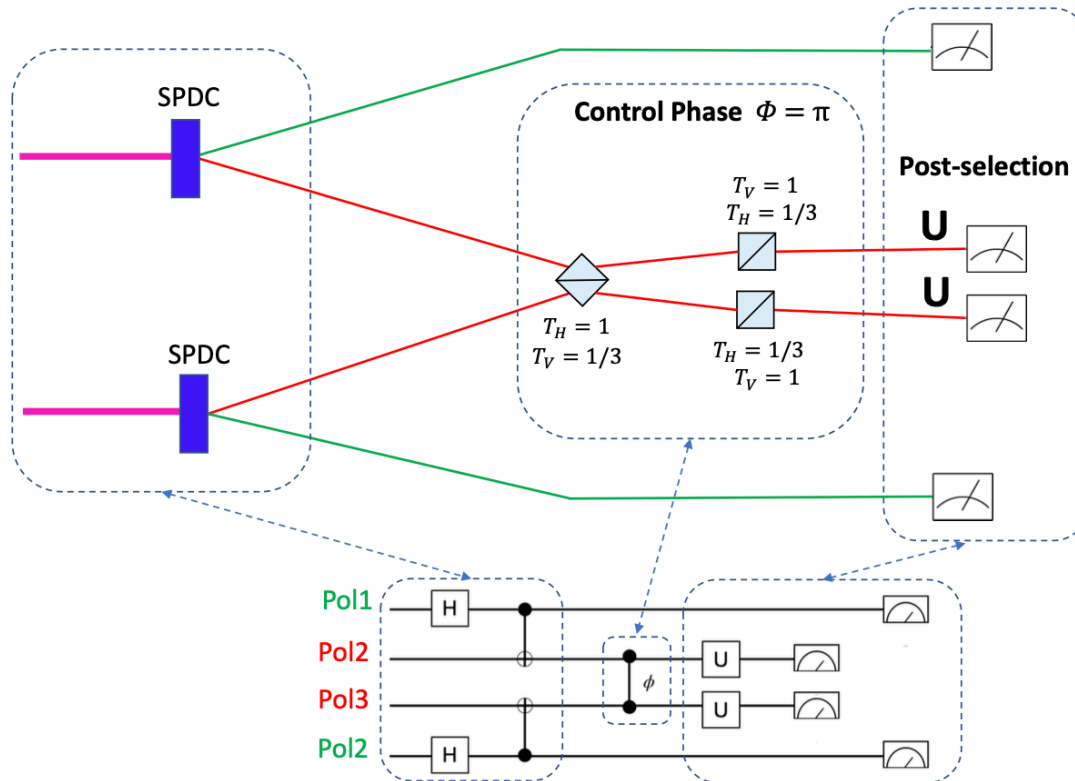


Optics Simulation



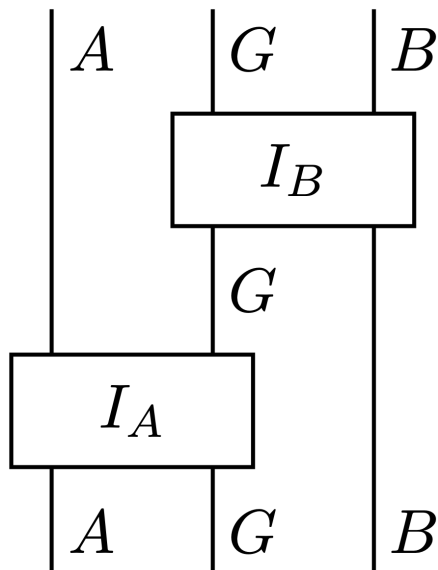
Optics Simulation

Four-photon scheme: Each qubit mapped to photon path



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Local Operations and Classical Communication cannot create entanglement.

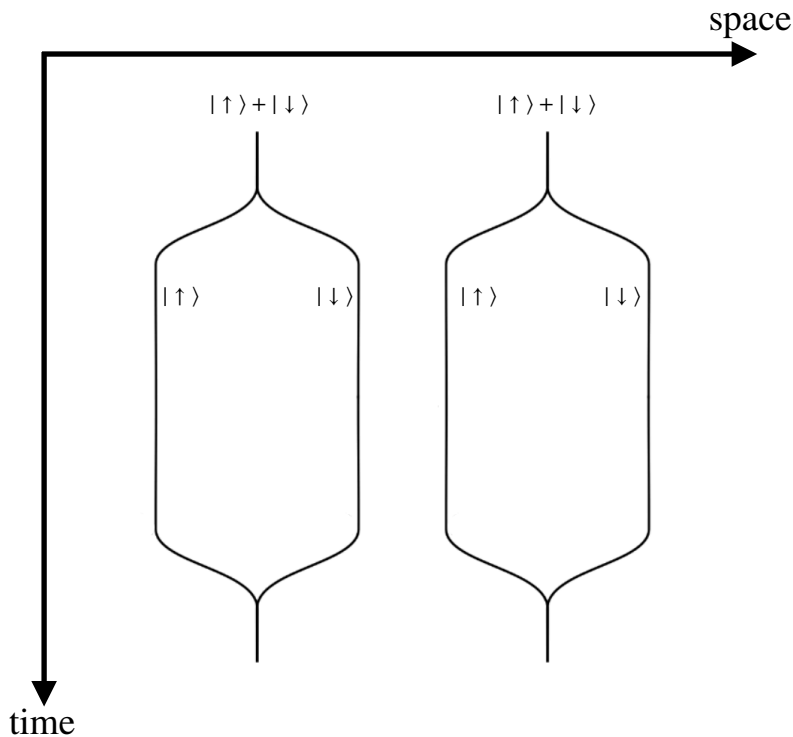


"If A and B get entangled by interacting only with G, then G cannot be classical."

Formalised in QI and GPTs.

The experiment can rule out a class of theories.

Instantaneous interaction?



Extant derivations of the effect made use of the static approximation.

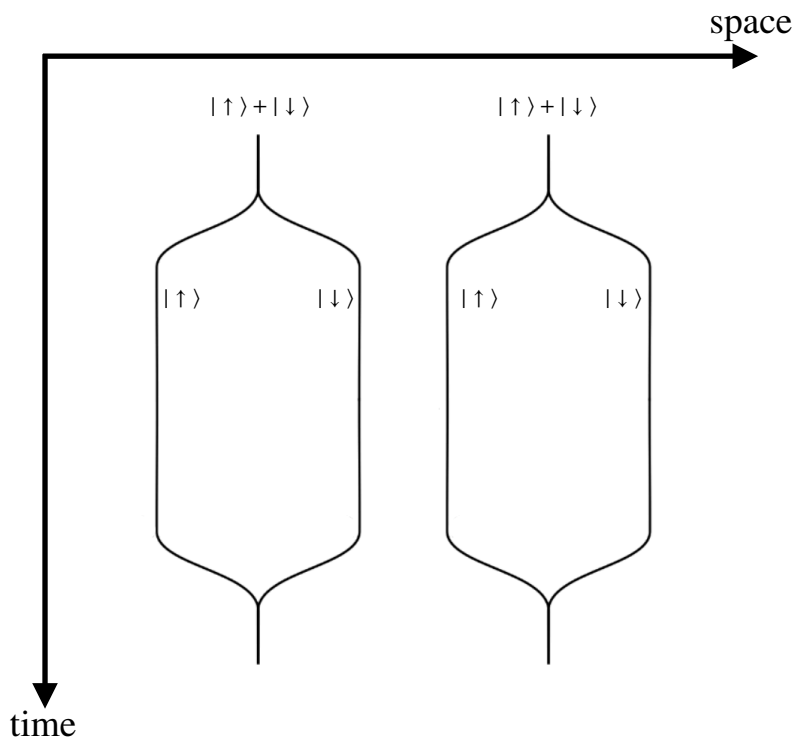
In that case, the effect can be explained by a direct interparticle interaction:

$$\hat{H}_{\text{int}} = - \frac{Gm_A m_B}{|\hat{x}_A - \hat{x}_B|}$$

We make use of the path-integral formulation of QM to compute the phases developed during the experiment.

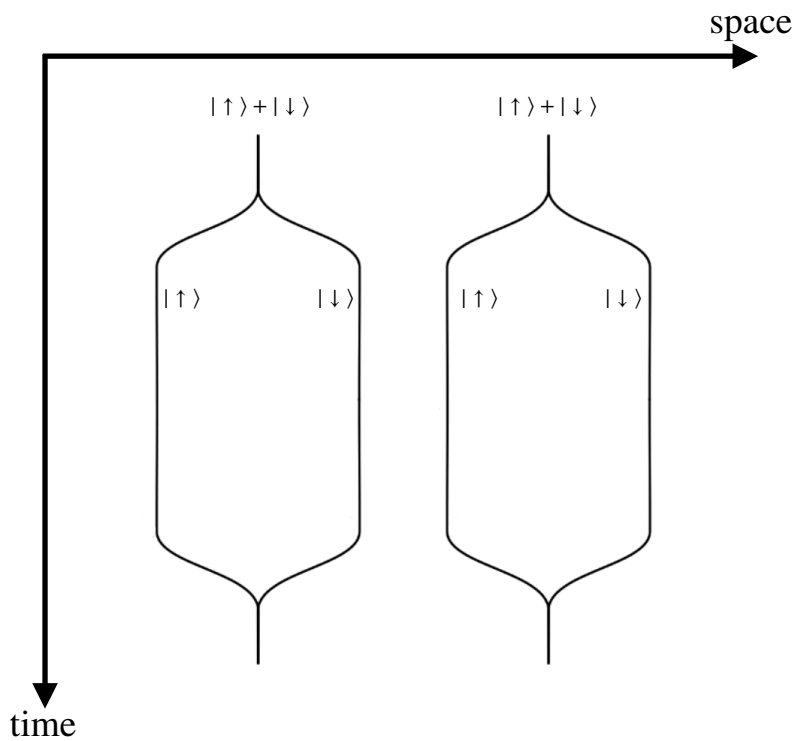
Path Integral

$$Z \propto \int \mathcal{D}x_1 \mathcal{D}x_2 \mathcal{D}F e^{iS[x_1, x_2, F[x_1, x_2]]/\hbar}$$



Path Integral

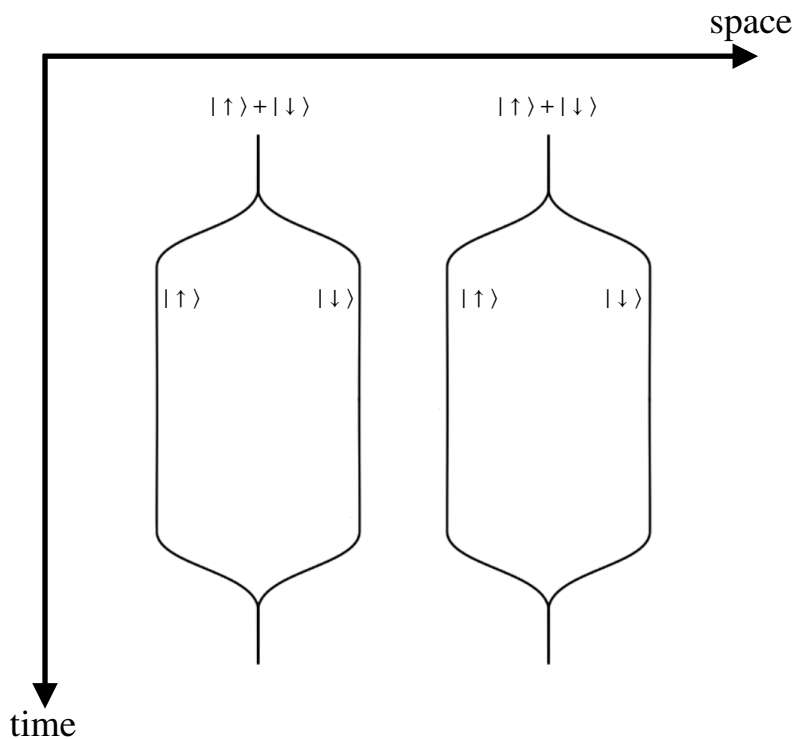
$$Z \propto \int \mathcal{D}x_1 \mathcal{D}x_2 \mathcal{D}F e^{iS[x_1, x_2, F[x_1, x_2]]/\hbar}$$



Assume the paths of the particles is imposed by the interaction with an external field.

Path Integral

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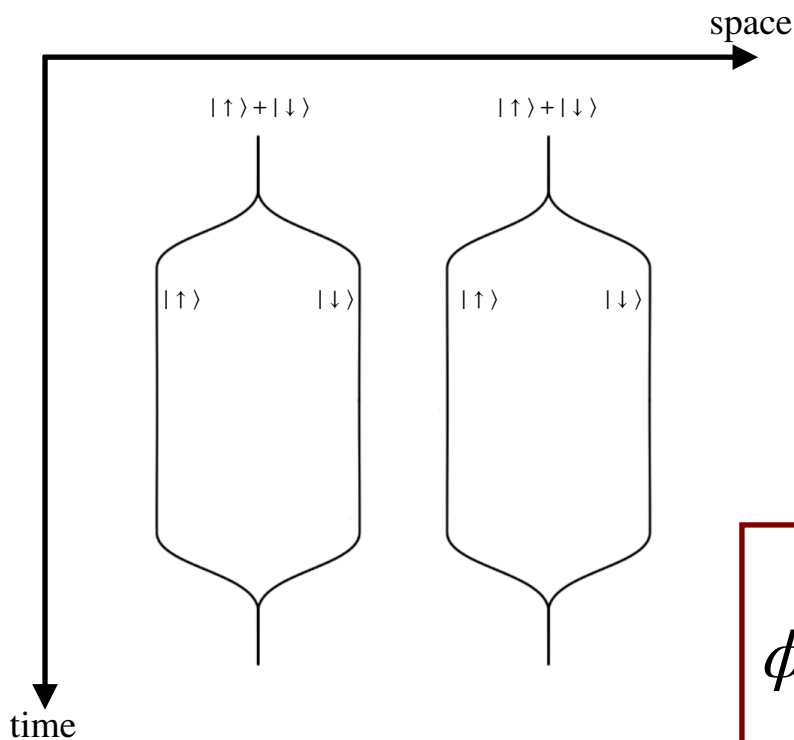


Assume the paths of the particles is imposed by the interaction with an external field.

Stationary phase approximation.

Path Integral

$$Z \propto \int \mathcal{D}x_1 \mathcal{D}x_2 \mathcal{D}F e^{iS[x_1, x_2, F[x_1, x_2]]/\hbar}$$



Assume the paths of the particles is imposed by the interaction with an external field.

Stationary phase approximation.

$$\phi(s_1, s_2) = \frac{iS^{\text{os}}[x_1^{s_1}, x_2^{s_2}, F[x_1^{s_1}, x_2^{s_2}]]}{\hbar}$$

$$\phi(s_1, s_2) = \frac{iS^{\text{OS}} [x_1^{s_1}, x_2^{s_2}, F[x_1^{s_1}, x_2^{s_2}]]}{\hbar}$$

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- Phases are actions, have the same symmetries

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- Phases are actions, have the same symmetries
- Gauge-invariant + Lorentz covariant (manifestly local)

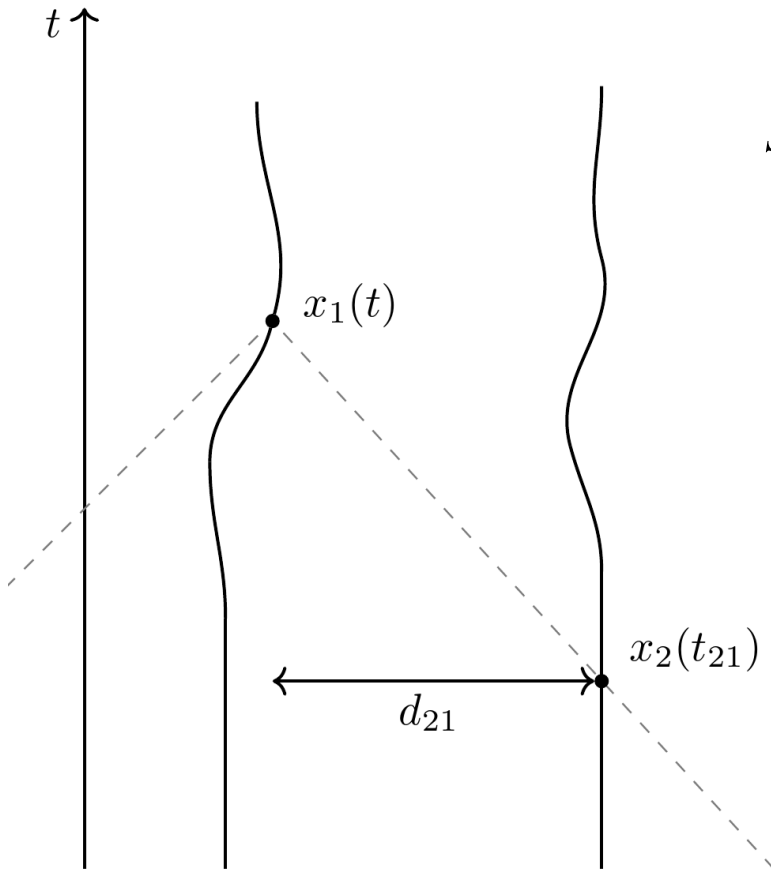
$$\phi(s_1, s_2) = \frac{iS^{\text{OS}}[x_1^{s_1}, x_2^{s_2}, F[x_1^{s_1}, x_2^{s_2}]]}{\hbar}$$

- Phases are actions, have the same symmetries
- Gauge-invariant + Lorentz covariant (manifestly local)
- Can be computed for arbitrary particle trajectories.

Gravitational phases

Exact formula

$$S_h = \frac{G}{c^4} \int dt \left(\frac{m_1 m_2 V_{1\mu\nu}(t) \bar{V}_2^{\mu\nu}(t_{21})}{|d_{21}(t)| - d_{21}(t) \cdot v_2(t_{21})/c} + 1 \leftrightarrow 2 \right)$$



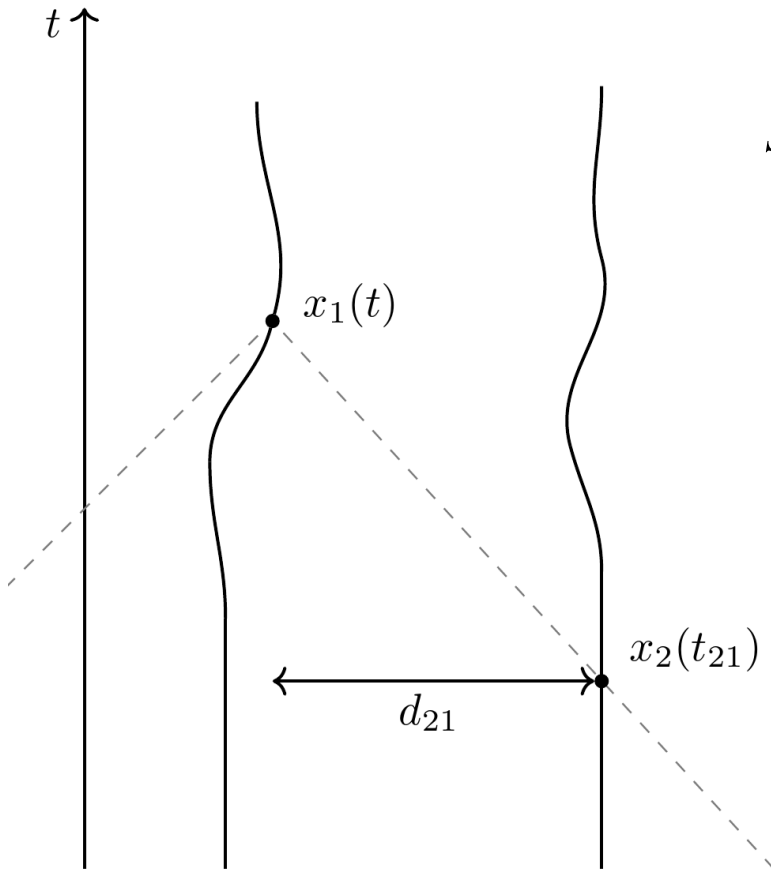
Gravitational phases

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Small velocities

$$S_h = \frac{G}{2} \int dt \left(\frac{m_1 m_2}{|d_{21}(t)|} + \frac{m_1 m_2}{|d_{12}(t)|} \right)$$



Gravitational phases

Exact formula

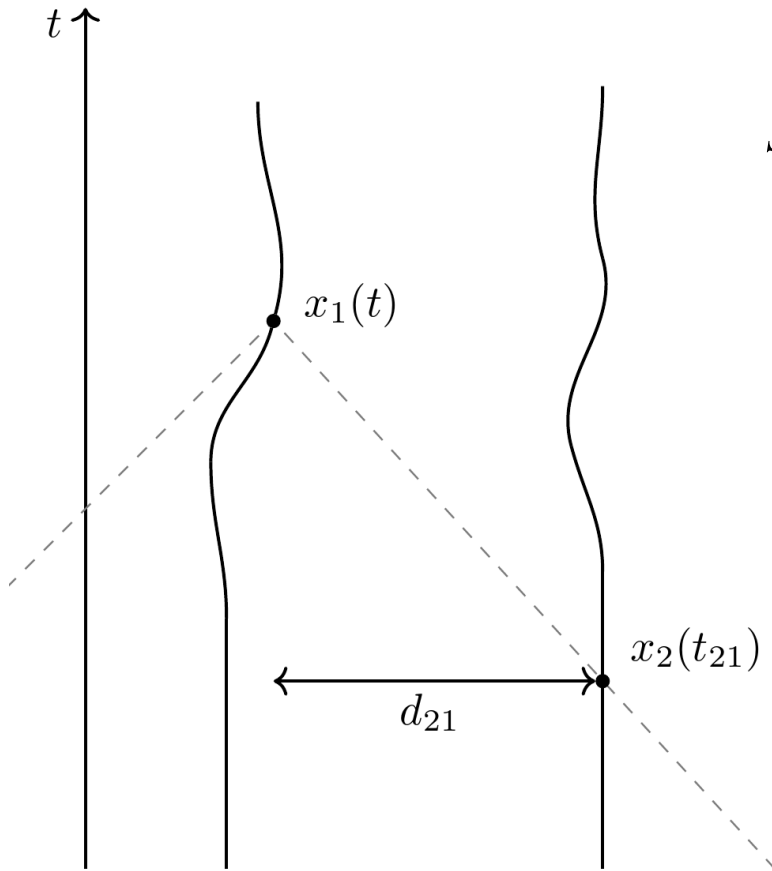
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Small velocities

$$S_h = \frac{G}{2} \int dt \left(\frac{m_1 m_2}{|d_{21}(t)|} + \frac{m_1 m_2}{|d_{12}(t)|} \right)$$

Newtonian

$$S_F = \int dt \frac{G m_1 m_2}{|d(t)|}$$



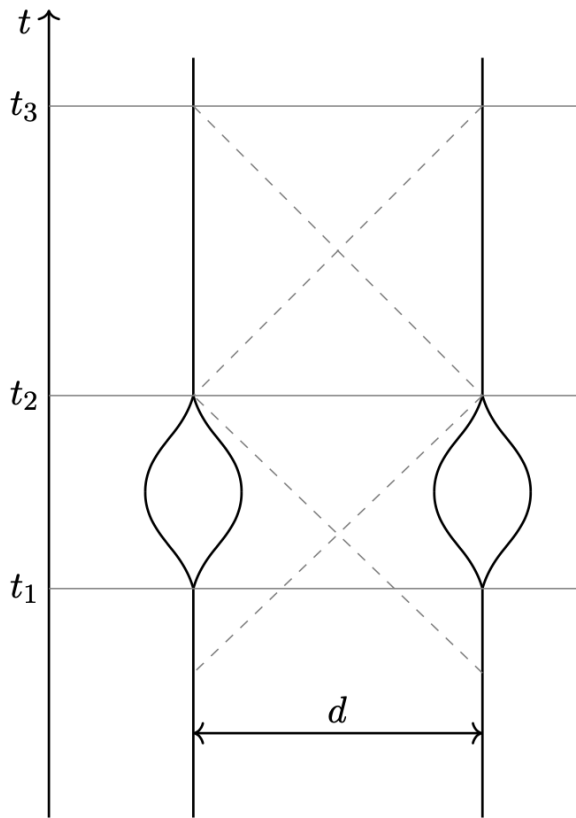
Observable effects

$$\phi(s_1, s_2) = \frac{G}{2\hbar} \int dt \left(\frac{m_1 m_2}{|d_{21}(t)|} + \frac{m_1 m_2}{|d_{12}(t)|} \right)$$

If superposition happens in spacelike separated regions \implies no entanglement!

Signal of the superposition needs time to propagate casually between masses.

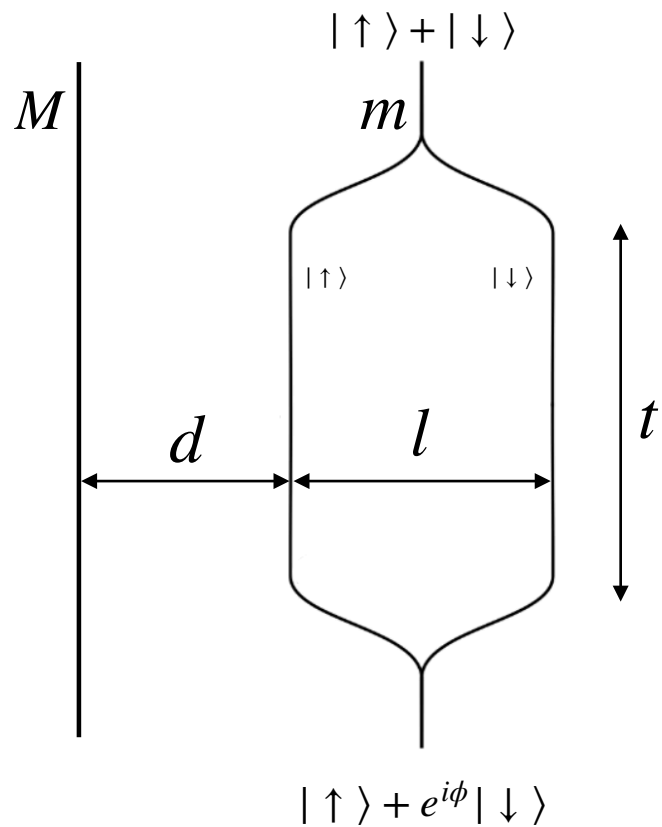
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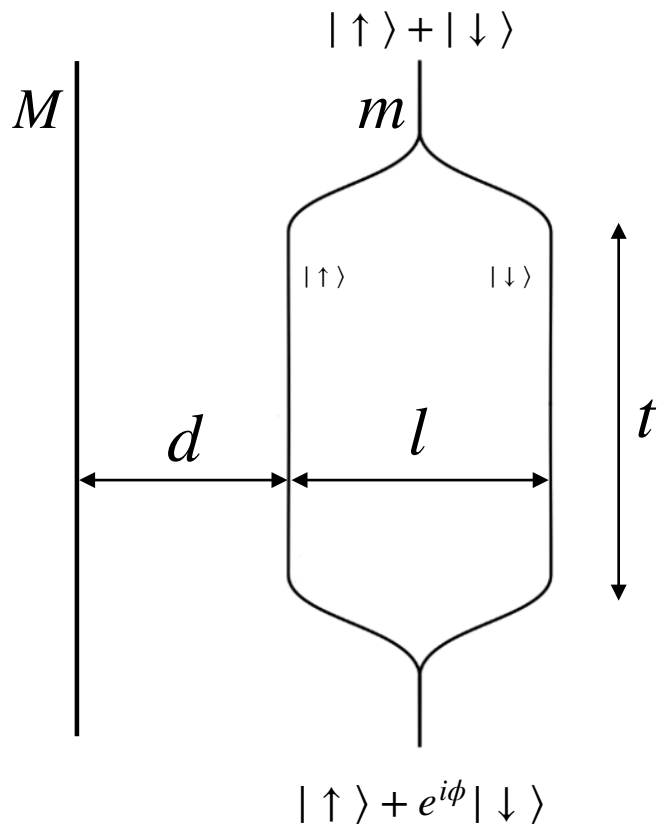
Smaller deviations due to retarded interaction could be measured in electron interferometry—for the EM case!

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Setup



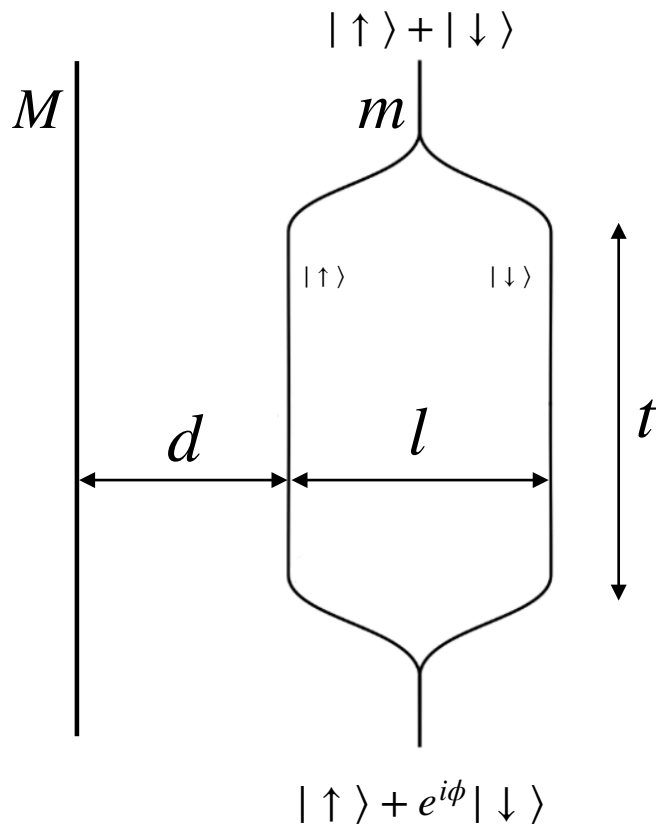
Setup



Time dilation

$$\tau(r) = \sqrt{|g_{00}(r)|} t = \sqrt{1 - \frac{2GM}{r}} t \approx \left(1 - \frac{GM}{r}\right) t$$

Setup



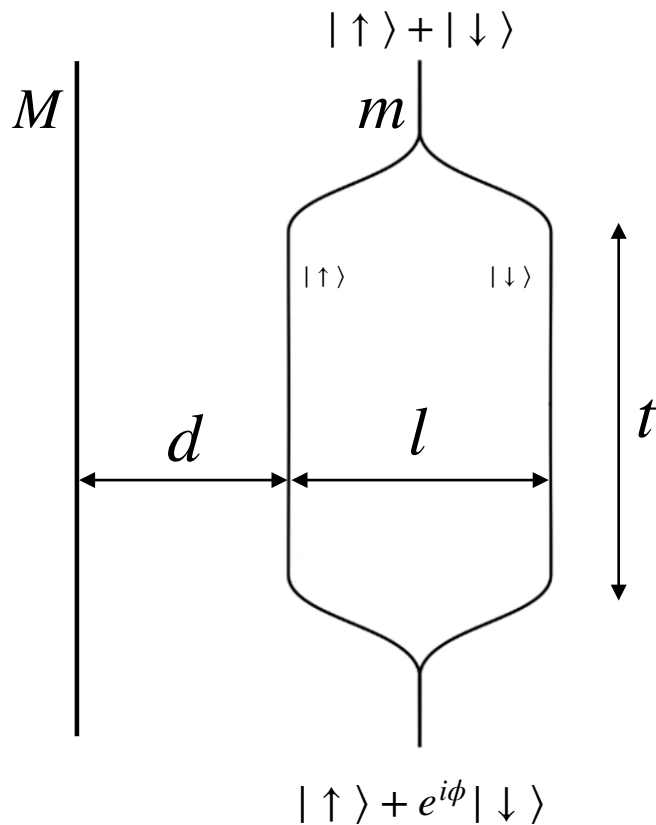
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Phases due to proper time differences

$$\phi = \frac{mc^2}{\hbar} \delta\tau = \frac{m}{m_P} \frac{\delta\tau}{t_P}$$

Setup



Time dilation

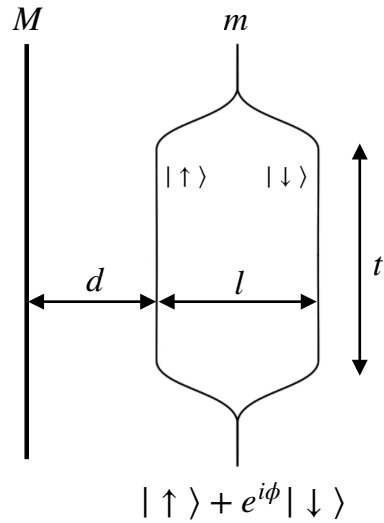
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Phases due to proper time differences

$$\phi = \frac{mc^2}{\hbar} \delta\tau = \frac{m}{m_P} \frac{\delta\tau}{t_P}$$

$$\delta\tau = \frac{GM}{c^2} \frac{l}{d(d+l)} t$$

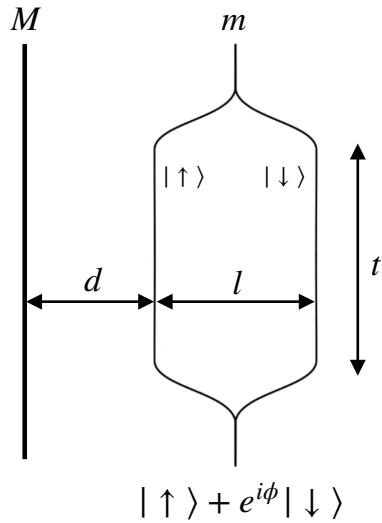
Hypothesis



$$\phi = \frac{m}{m_P} \frac{\delta\tau}{t_P}$$

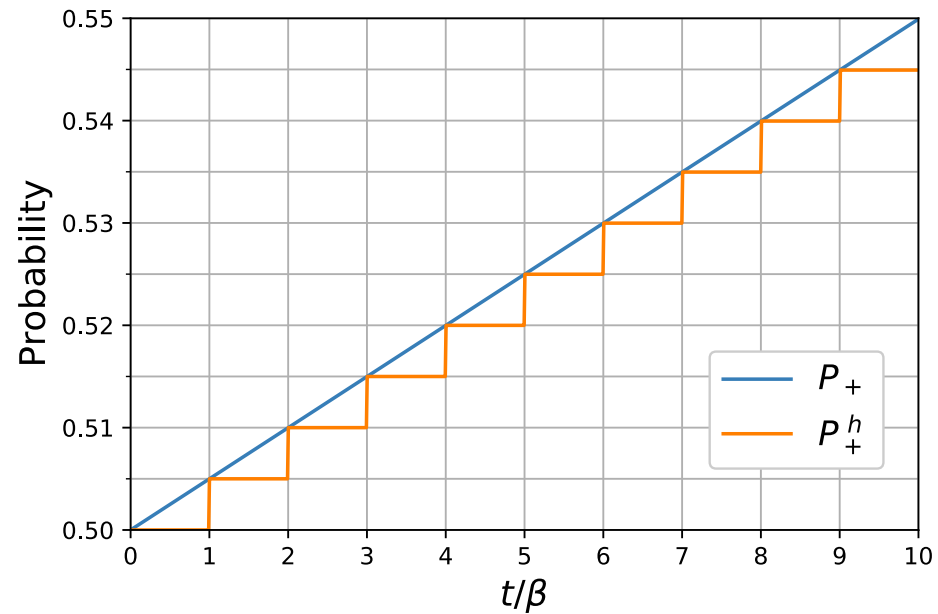
$$\delta\tau = nt_P, \quad n \in \mathbb{Z}$$

Hypothesis



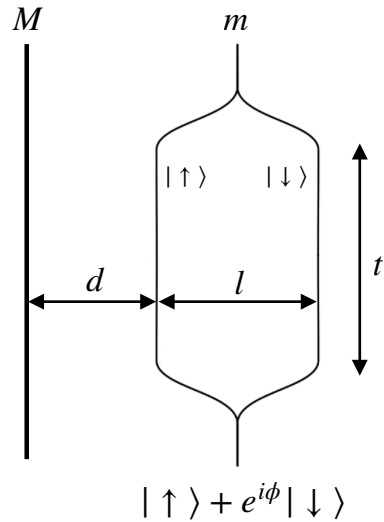
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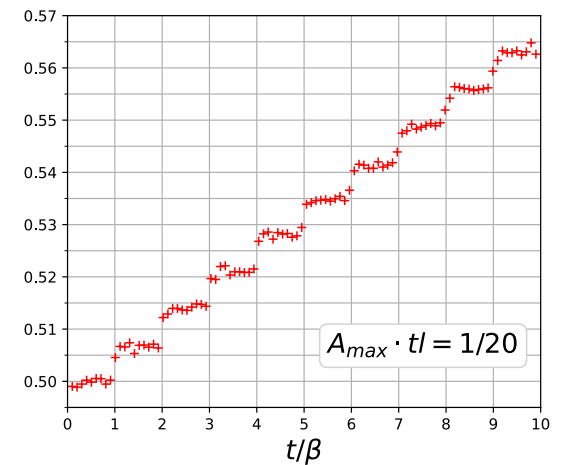
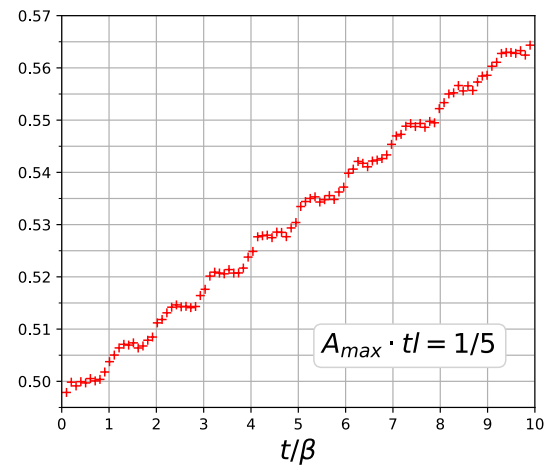
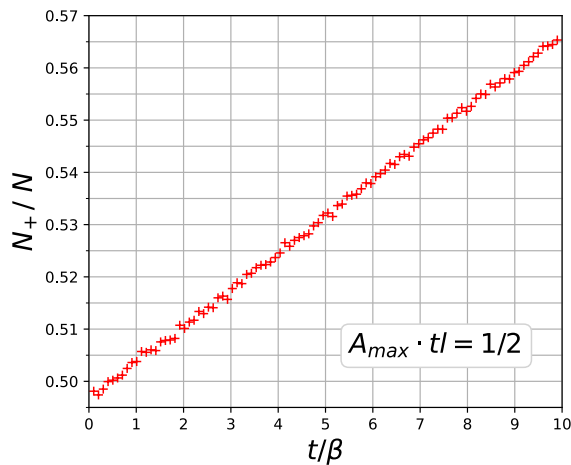
$$\delta\tau = \left\lfloor \frac{t}{\beta} \right\rfloor t_P$$

Results



Parameter	Value	Uncertainty
m	3×10^{-10} kg	10^{-12} kg
M	3×10^{-9} kg	10^{-11} kg
t	10^{-1} s	10^{-4} s
l	10^{-7} m	10^{-9} m
d	[17, 54] cm	10^{-2} cm
A	$\leq 4 \times 10^{-10}$ kg m $^{-2}$	

$$T \approx 4 \text{ K} \quad P \approx 10^{-18} \text{ Pa}$$



- **Part I: Quantum gravity (and beyond) in the lab**
- **Part II: Conceptual investigations**

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 - The arrow of time in operational formulations of QT
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Tensions

Operational formulations of QM are strongly time-oriented.

Quantum states are associated with the past of a system.

Probabilities are about future results.



In tension with time-reversal symmetry of the rest of fundamental physics.

An issue for the reconstructions of quantum mechanics.

Does quantum uncertainty imply time orientation?

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No.

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Then why are certain formulations of quantum theory time-oriented?

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They are designed to describe the interaction of macroscopic thermodynamical systems with quantum systems.

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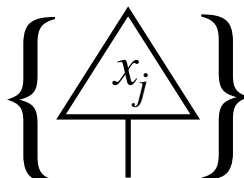
Then why are certain formulations of quantum theory time-oriented?

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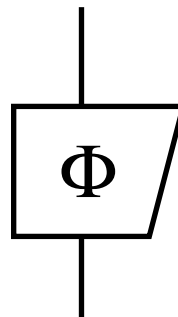
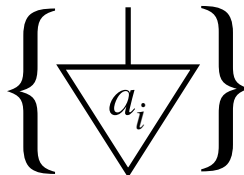
We remember the past, but not the future.

Two games

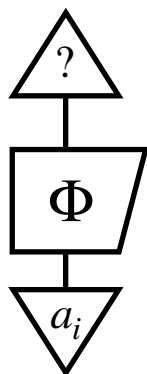
Measurement



Preparation



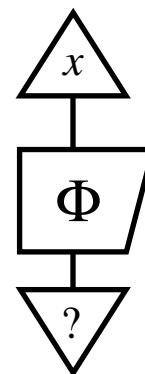
Evolution



$$P_{pre}(x | a, \Phi)$$

Prediction

$$P_{post}(a | x, \Phi)$$



Postdiction

Inference Symmetry

A process Φ is **inference symmetric** if:

$$P_{pre}(x_j | a_i, \Phi) = P_{post}(a_i | x_j, \Phi)$$

for any choice of bases.

A kind of passive time-reversal symmetry.

Unitary evolution is inference symmetric.

Quantum channels are not inference symmetric.

Inference Symmetry

A process Φ is **inference symmetric** if:

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A kind of passive time-reversal symmetry.

Unitary evolution is inference symmetric.

Quantum channels are not inference symmetric.

The inference asymmetry of quantum channels is understood as an asymmetry in the inference data.

$$P_{pre}(x | a, \Phi) = \begin{array}{c} \triangle x \\ | \\ \square \Phi \\ | \\ \triangle a \end{array} = \begin{array}{c} \triangle x \quad \equiv \\ | \quad | \\ \square U_{\Phi} \\ | \quad | \\ \triangle a \quad \triangle b \end{array} = P_{pre}(x | ab, U_{\Phi})$$

$$P_{post}(a | x, \Phi) = P_{post}(a | xb, U_{\Phi})$$

Why the asymmetry?

Time-asymmetry due to the users of QM.

QI is about correlations established between agents.

The agent is not explicitly modelled by the theory, but *represented* in the mathematical objects in the theory.

- **Part II: Conceptual investigations**
 - The arrow of time in operational formulations of QT
 - **The relational interpretation of QM.**

Map of Madness

	ψ -Ontic	ψ -Epistemic
Type-I (intrinsic realism)	Bohmian mechanics ^{10,11} Many worlds ^{12,13} Modal ^{14,15} Bell's "beables" ¹⁶ Collapse theories* ^{17,18}	Einstein ¹⁹ Ballentine ²⁰ Consistent histories ^{21,22} Spekkens ²³

	About knowledge	About belief
Type-II (participatory realism)	Copenhagen ^{24,25} Wheeler ^{26,27} Relational ^{28,29} Zeilinger ^{3,30} No "interpretation" ³¹ Brukner ³²	QBism ³³⁻³⁵

+ objective collapse models: Penrose-Diósi, GRW...

arXiv:1509.04711 (quant-ph)

[Submitted on 15 Sep 2015 (v1), last revised 23 Nov 2016 (this version, v2)]

Interpretations of quantum theory: A map of madness

Adán Cabello

Interpretations of quantum mechanics:

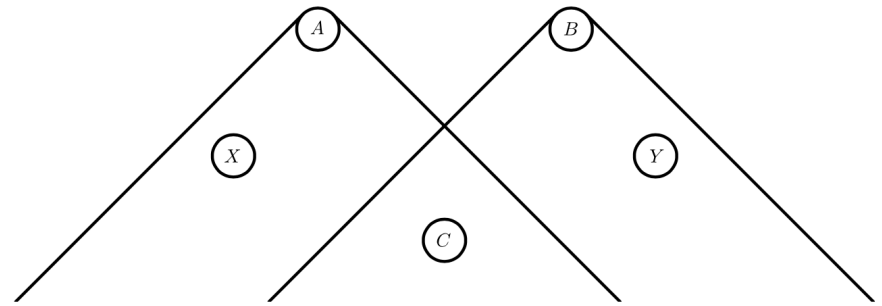
- Surprisingly different pictures of the world
- Designed to give the same predictions (except for objective collapse)
- But experimental metaphysics *can* put constraints on them.

No-Go theorems

Put constraints on various features of an interpretation.

Bell's 1967 theorem says QM is incompatible with:

- Relativistic causality
- Reichenbach's principle of decorrelating explanation
- No Superdeterminism



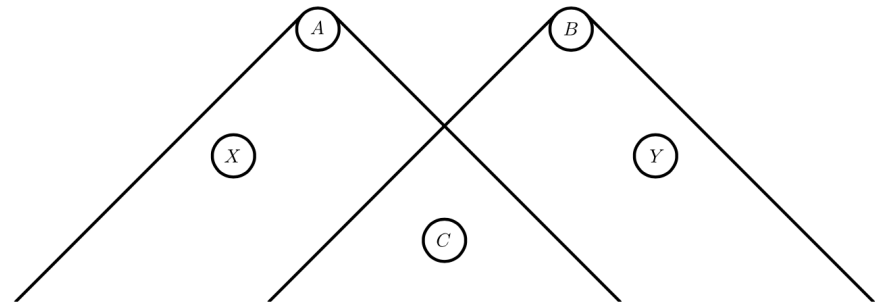
$$P(ab|c) = P(a|c)P(b|c)$$

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$$P(ab|c) = P(a|c)P(b|c)$$

Implicit assumption: Absoluteness of observed events

No-Go theorems

Recent theorem by Bong *et.al.* shows that QM is incompatible with

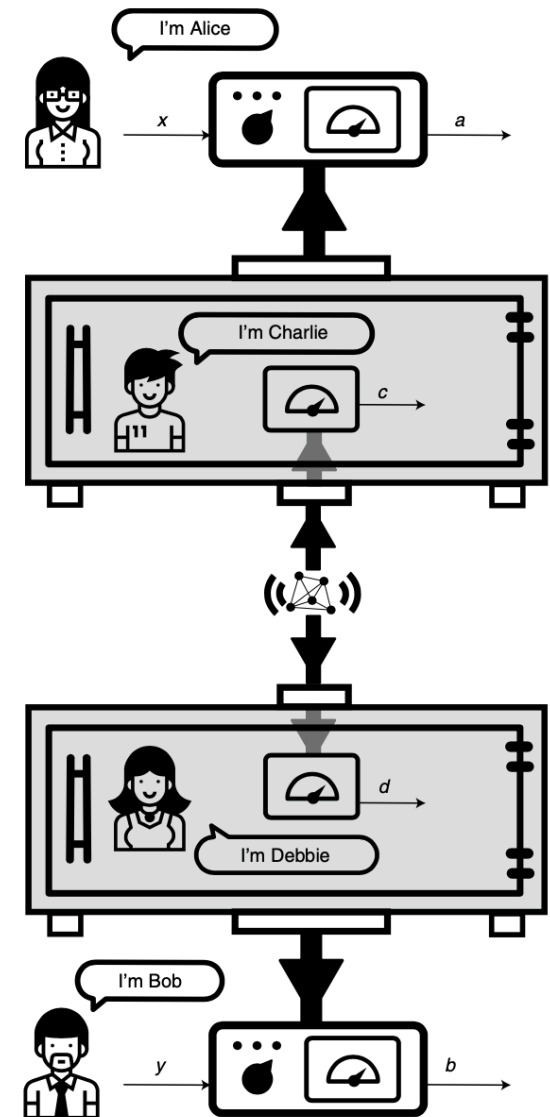
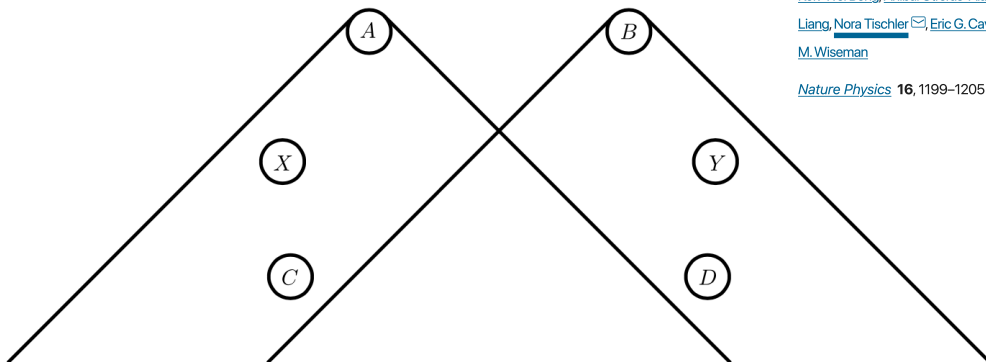
- Locality
- No Superdeterminism
- **Absoluteness of observed events**

Article | [Published: 17 August 2020](#)

A strong no-go theorem on the Wigner's friend paradox

[Kok-Wei Bong](#), [Anibal Utreras-Alarcón](#), [Farzad Ghafari](#), [Yeong-Cherng Liang](#), [Nora Tischler](#) ✉, [Eric G. Cavalcanti](#) ✉, [Geoff J. Pryde](#) & [Howard M. Wiseman](#)

Nature Physics **16**, 1199–1205 (2020) | [Cite this article](#)



Relational Quantum Mechanics is an interpretation of QM that embraces the relativity of facts.

In RQM, facts are relations established between two systems.

What is a fact relative to a given system might not be a fact relative to another.

"Wigner's facts are not necessarily his friend's facts"

Facts can happen relative to *any* physical system.

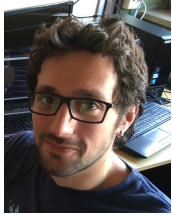
How does the classical world emerge from the world of relative facts?

To what extent the relativity of facts is analogous with special relativity?

How can objectivity be achieved when facts are not shared?

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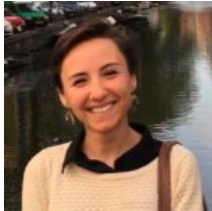
Collaborators



**Davide
Poderini**



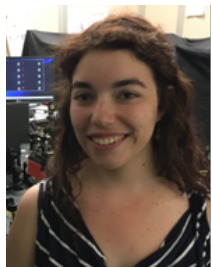
**Emanuele
Polino**



**Beatrice
Polacchi**



**Fabio
Sciarrino**

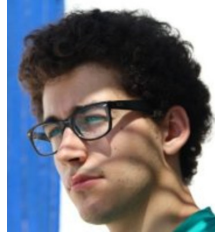


**Iris
Agresti**



**Gonzalo
Carvacho**

**Photonic Implementation of
Quantum Gravity Simulators**



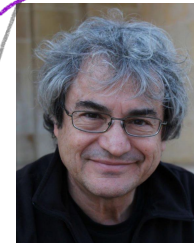
**Pierre
Martin-Dussaud**

[Submitted on 16 Jul 2020 (v1), last revised 31 Jan 2022 (this version, v4)]

**An experiment to test the
discreteness of time** 2007.08431



**Marios
Christodoulou**



**Carlo
Rovelli**



Pietro Donà

[Open Access](#) | [Published: 27 February 2021](#)

Stable Facts, Relative Facts

[Foundations of Physics](#) 51, Article number: 30 (2021)

[Submitted on 7 Oct 2021] 2202.03368

**Relational Quantum Mechanics is
about Facts, not States: A reply to
Pienaar and Brukner**



**Markus
Aspelmeyer**



**Časlav
Brukner**



**Richard
Howl**

[Submitted on 7 Feb 2022] 2202.03368

**Locally mediated entanglement
through gravity from first principles**

The arrow of time in operational
formulations of quantum theory

2021-08-09, volume 5, page 520

[arXiv:2010.05734v2](#)

<https://doi.org/10.22331/q-2021-08-09-520>

**Thank
you!**